

# Peculiar architectures for the WASP-53 and WASP-81 planet-hosting systems<sup>★</sup>

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## ABSTRACT

We report the detection of two new systems containing transiting planets. Both were identified by WASP as worthy transiting planet candidates. Radial velocity observations quickly verified that the photometric signals were indeed produced by two transiting hot Jupiters. Our observations also show the presence of additional Doppler signals. In addition to short-period hot Jupiters, we find that the WASP-53 and WASP-81 systems also host brown dwarfs, on fairly eccentric orbits with semimajor axes of a few astronomical units. WASP-53c is over  $16 M_{\text{Jup}} \sin i_c$  and WASP-81c is  $57 M_{\text{Jup}} \sin i_c$ . The presence of these tight, massive companions restricts theories of how the inner planets were assembled. We propose two alternative interpretations: the formation of the hot Jupiters within the snow line or the late dynamical arrival of the brown dwarfs after disc dispersal. We also attempted to measure the Rossiter–McLaughlin effect for both hot Jupiters. In the case of WASP-81b, we fail to detect a signal. For WASP-53b, we find that the planet is aligned with respect to the stellar spin axis. In addition we explore the prospect of transit-timing variations, and of using *Gaia*'s astrometry to measure the true masses of both brown dwarfs and also their relative inclination with respect to the inner transiting hot Jupiters.

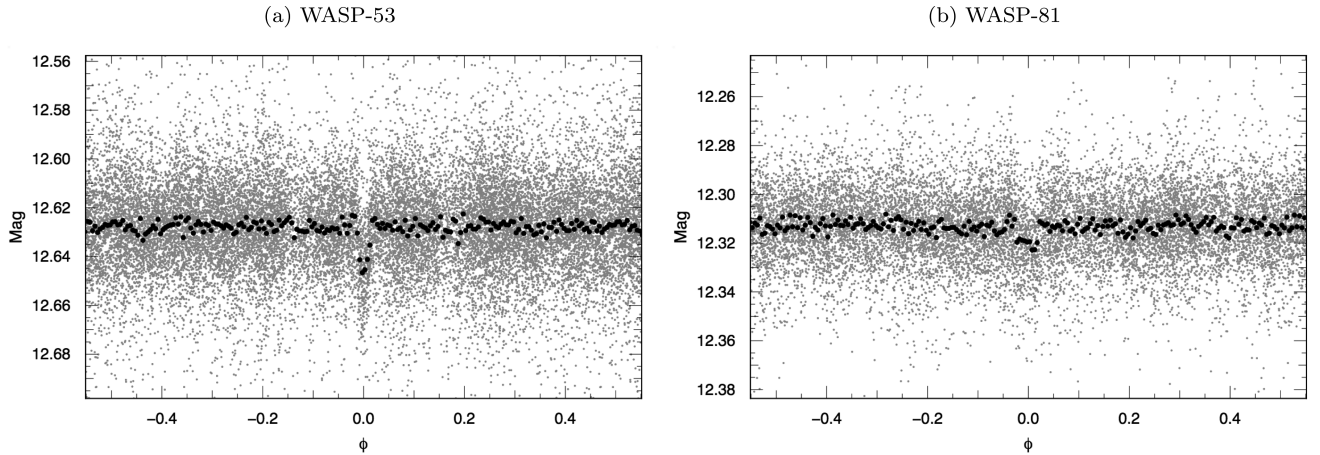
**Key words:** planets and satellites: individual: WASP-81 – planets and satellites: individual: WASP-53 – binaries: eclipsing – brown dwarfs – planetary systems.

## 1 FOREWORDS

The discovery of 51 Peg b (Mayor & Queloz 1995) initiated a debate about the origin (formation and evolution) of hot Jupiters that continues to rage to this day. The question is whether they formed *in situ*, within the so-called snow line (Bodenheimer, Hubickyj & Lissauer 2000; Lee & Chiang 2015; Batygin, Bodenheimer & Laughlin 2016), or beyond, followed by inward migration (Pollack et al. 1996; Alibert et al. 2005; Helled et al. 2014). The second

<sup>★</sup> using data collected at ESO's La Silla Observatory, Chile: HARPS on the ESO 3.6m (Prog IDs 087.C-0649, 089.C-0151, 090.C-0540, 091.C-0184 & 093.C-0474), the ESO NTT (Prog ID 088.C-0204), the Swiss *Euler* telescope, and TRAPPIST. The data is publicly available online, and on demand to the main author.

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**Figure 1.** Variations in the magnitude of WASP-53 and WASP-81 leading to their identification as transiting planet candidates. The grey dots show individual WASP measurements, whereas the black dots show the median magnitude within each of 200 phase bins.

hypothesis needs an explanation of how hot Jupiters have migrated, either via angular momentum transfer with their protoplanetary disc (Lin, Bodenheimer & Richardson 1996; Ward 1997; Baruteau et al. 2014) or thanks to dynamical interactions followed by tidal circularization (high-eccentricity migration; Rasio & Ford 1996; Wu, Murray & Ramsahai 2007; Naoz et al. 2011; Petrovich 2015a). Any framework has to explain why gas giants are found both close and far from their host star (not least Jupiter and Saturn), and also that they frequently orbit on planes that are inclined, sometimes retrograde, with respect to the equatorial plane of their host star (Hébrard et al. 2008; Winn et al. 2009; Anderson et al. 2010; Schlaufman 2010; Triaud et al. 2010; Albrecht et al. 2012a; Lendl et al. 2014; Winn & Fabrycky 2015).

Observations of the spin-orbit angle, thanks to the Rossiter-McLaughlin effect (Queloz et al. 2000), were initially thought to provide a clean test between disc-driven migration and dynamical and tidal migration (Gaudi & Winn 2007). While some studies indicate that observations are compatible with planets undergoing orbital realignment (Triaud 2011; Albrecht et al. 2012b; Dawson 2014), additional theoretical arguments imply that misalignments can arise via a variety of processes. They can be primordial, with the planets forming on inclined planes (Lai, Foucart & Lin 2011; Thies et al. 2011; Batygin 2012; Lai 2014; Spalding & Batygin 2015), or they can arise later (Cébron et al. 2011; Rogers et al. 2013). Thus inclined hot Jupiters are compatible with both dynamical interactions and with disc-driven migration.

Obviously, disc-driven migration and high-eccentricity migration could both be correct and each produce a fraction of the hot Jupiters, as recent observation appear to imply (Anderson et al. 2015). Finally, results from Guillochon, Ramirez-Ruiz & Lin (2011) suggest that high-eccentricity migration requires in most cases some amount of disc-driven migration.

In this paper, we spend the first sections to describe the discovery of two new planetary systems. Both present a tight hierarchical architecture, composed of an inner, transiting, hot Jupiter, and an outer brown dwarf companion. WASP-53 and WASP-81 are reminiscent of a challenge proposed by Hatzes & Wuchterl (2005) to test disc-driven migration. If both brown dwarfs were on their current orbits during the protoplanetary phase, they would have truncated the disc within the *snow line* thus preventing disc-driven migration and only allowing in situ formation.

In the final sections we describe a number of formation and evolution scenarios, some involving disc-driven migration and others not.

We conclude that a number of scenarios will become testable soon, with the arrival of precise astrometric measurements produced by ESA's *Gaia* satellite. We also compute the equilibrium eccentricities of the inner gas giants (Mardling 2007) and find it is unlikely that their  $k_2$  Love number (Batygin, Bodenheimer & Laughlin 2009) will be measured soon.

## 2 WASP IDENTIFICATION

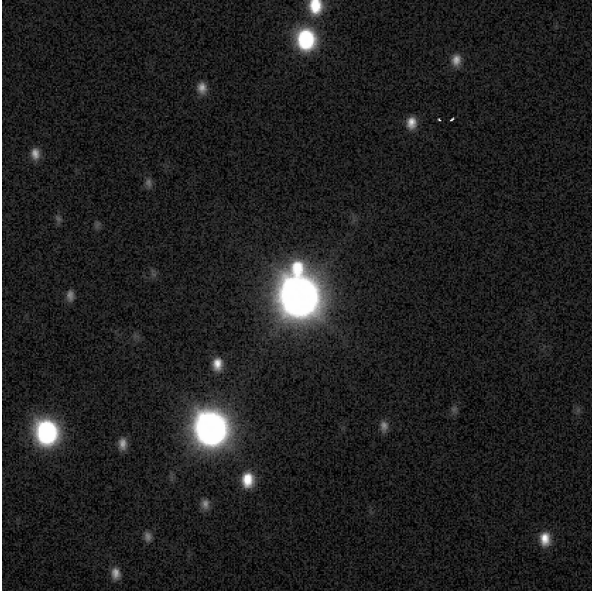
The WASP survey<sup>1</sup> (Pollacco et al. 2006) consists of two sets of eight 11-cm refractive telescope mounted together. One set is located at the Observatorio del Roque de Los Muchachos, La Palma (Spain), while the other is installed in Sutherland, hosted by the South African Astronomical Observatory. WASP has observed in excess of 30 million stars since 2004, thousands of times each. The photometric data reduction and the candidate selection are described in Collier Cameron et al. (2007).

WASP-53 (2MASS J02073820–2039426; K3,  $J = 10.959$ ) and WASP-81 (2MASS J20164989+0317385; F9,  $J = 11.263$ ) have been observed 21 120 and 13 292 times by WASP. They were two unremarkable and anonymous stars before short, box-like photometric signals were identified at  $P = 3.309866$  and  $2.716554$  d, respectively. WASP-53b was sent for radial velocity verification in 2010-08-10, and WASP-81b in 2011-05-09, with the first spectra acquired on 2010 Decemehr 5 and 2011 September 29. The WASP data are shown in Fig. 1. We temporarily call a visual companion to WASP-53A, WASP-53B, which is show in Fig. 2, but find the pair is likely unrelated (see Section 3.1.1).

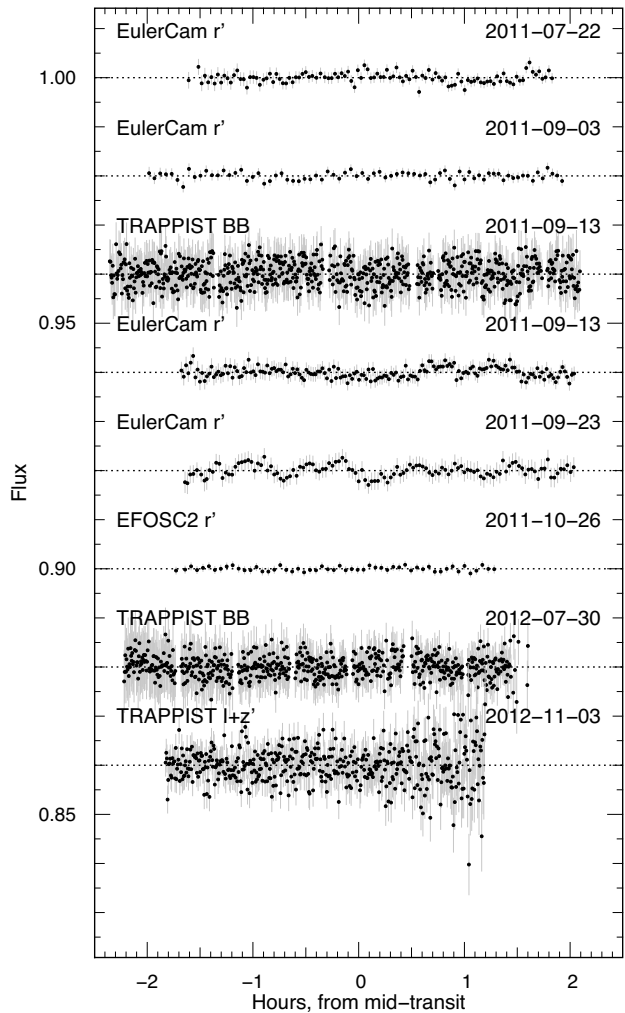
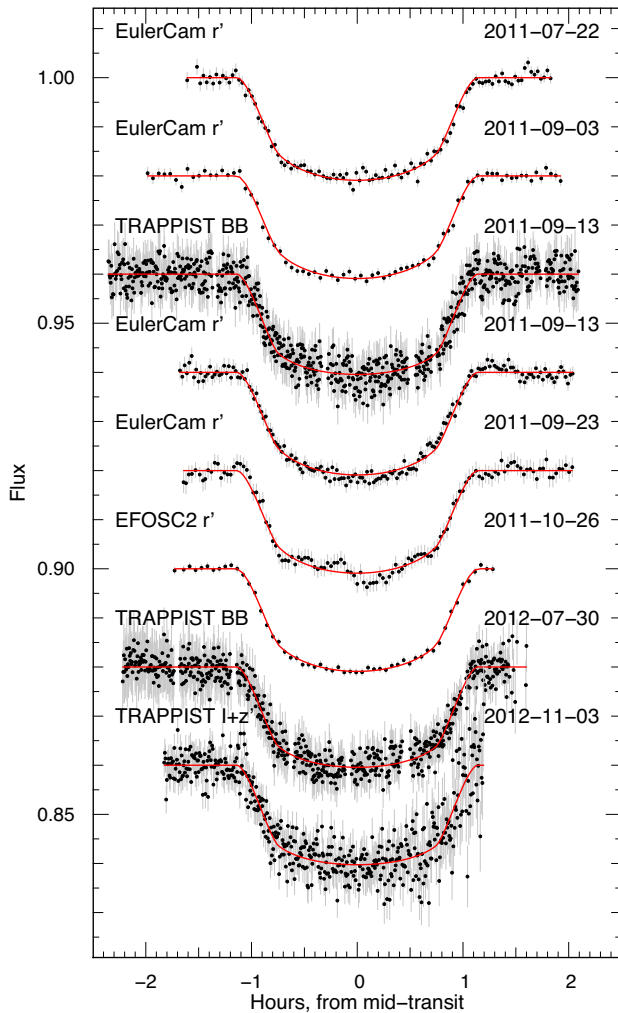
## 3 PHOTOMETRIC OBSERVATIONS

All follow-up photometric time series, their dates, filters, number of data points and the detrending functions that were used in the analysis are detailed in Table C1. Fig. 3 shows the corrected photometry for WASP-53, and Fig. 4 shows WASP-81. Raw fluxes and residuals are shown in Figs C1 and C2, respectively. Data used in the analysis are presented in Appendix A and available online. Below, we provide some details on the observations and reduction, although we encourage readers to refer to cited papers for full information.

<sup>1</sup> <https://wasp-planets.net>



**Figure 2.** A  $1.5 \times 1.5$  arcmin<sup>2</sup> field of view centred on WASP-81, obtained with EulerCam, clearly showing a visual companion to the north. The pair is most likely unrelated.



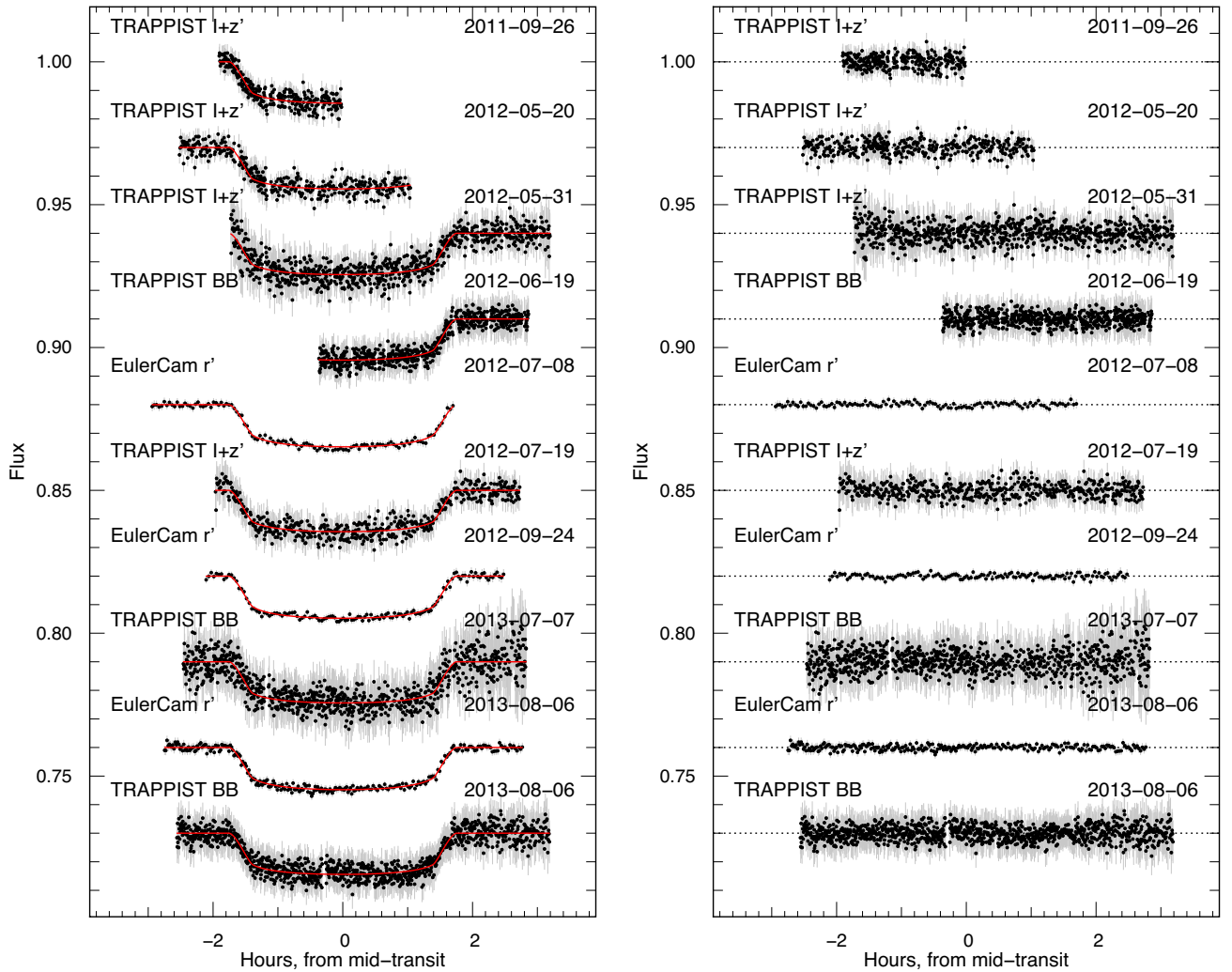
**Figure 3.** Photometry on WASP-53 at the time of transit, using EulerCam, TRAPPIST and the NTT. Left: we have the detrended data with, in red, the most likely model. Right: we show the residuals. Raw photometry and full models are available in Fig. C1.

### 3.1 EulerCam

EulerCam is mounted at the 1.2-m *Euler* Swiss telescope located at ESO La Silla Observatory (Chile). It has a pixel scale of 0.215 arcsec for a field of view of  $14.7 \times 14.5$  arcmin<sup>2</sup>. The telescope is an alt-azimuthal design. EulerCam is mounted behind a field de-rotator. To ensure the best photometric precision, each star is kept on the same pixels, using a digital feedback scheme that compares the newly acquired frame with a composite of earlier frames and their offset from a recorded position.

We obtained four transits of WASP-53b, all using an  $r'$ -Gunn filter, and with a slightly defocussed telescope to improve the observing efficiency and point spread function (PSF) sampling. Two of the transits, on 2011 September 22 and 2012 December 2, were scheduled to coincide with the radial velocity time series obtained using HARPS (see Section 4.2). WASP-81 was observed with EulerCam throughout three transits, one of which, on 2013 August 5, was obtained while HARPS collected a radial velocity time series. Our observations on 2012 June 7 were obtained using an  $I$ -Cousins filter and a focused telescope. The telescope was defocussed slightly for latter two observations, for which a  $z'$ -Gunn (2012-09-24) and an  $r'$ -Gunn (2013-08-05) were used.





**Figure 4.** Photometry on WASP-81 at the time of transit, using EulerCam and TRAPPIST. Left: we show the detrended data with, in red, the most likely model. Right: we show the residuals. Raw photometry and full models are available in Fig. C2.

All EulerCam images were reduced using standard image correction methods and light curves were obtained using differential aperture photometry with a careful selection of aperture and reference stars. For further details on the EulerCam instrument and the associated data reduction procedures, please refer to Lendl et al. (2012). The times of observations are provided in JD(UTC), changed later to BJD(TDB) during the global analysis.

### 3.1.1 A nearby stellar source

During spectroscopic observations, an additional light source was noticed near WASP-81A on CORALIE’s guiding camera images. On 2014 April 29, we obtained focused images of the WASP-81 system with EulerCam through Geneva *B* (six images), Geneva *V* (three images), and *r'*-Gunn (three images) filters (Fig. 2). We used astronomy.net (Lang et al. 2010) to calculate a precise astrometric solution and performed PSF fitting on the images using DAOPHOT (Stetson 1987). The results are presented in Table 1.

We extracted magnitudes for the visual companion relative to the primary target. Combining with apparent magnitudes for WASP-81A (Table 2), and using a  $E(B - V) = 0.05$  (Section 4), we obtained apparent magnitudes for the visual companion. We find

**Table 1.** Observation parameters for the companion source to WASP-81A.

Filter	$\Delta$ Mag	Mag	Separation (arcsec)	PA ( $^{\circ}$ )
<i>BG</i>	$5.64 \pm 0.03$	$18.78 \pm 0.07$	$4.34 \pm 0.02$	$3.5 \pm 0.2$
<i>VG</i>	$5.16 \pm 0.01$	$17.61 \pm 0.02$	$4.33 \pm 0.01$	$3.5 \pm 0.3$
<i>r'</i>	$4.87 \pm 0.01$	$17.14 \pm 0.04$	$4.32 \pm 0.01$	$3.5 \pm 0.3$
Weighted mean:			$4.33 \pm 0.01$	$3.5 \pm 0.3$

that the companion has colours consistent with a K3–K4 spectral type. If it is a dwarf, this implies a distance modulus of order 10.2 (1.1 kpc), compared to 8.0 (400 pc) for WASP-81A; if it were a giant it would be further away still. If we place the companion on a main-sequence isochrone (Marigo et al. 2008), and if WASP-81A were at the same distance, WASP-81A would have to be a 2-Gyr old,  $1.4\text{--}1.6 M_{\odot}$  star, contradicting our spectroscopic analysis as well as the mean stellar density obtained from the transit. If, instead, WASP-81A is on the main sequence, then the companion has to be below the main sequence to be at the same distance. The two objects are therefore most likely unrelated. We placed an HR diagram of the pair into the appendices.

**Table 2.** Stellar parameters of WASP-53A and WASP-81A.

Parameter	WASP-53A	WASP-81A
$\alpha$	02 <sup>h</sup> 07 <sup>m</sup> 38 <sup>s</sup> .22	20 <sup>h</sup> 16 <sup>m</sup> 49 <sup>s</sup> .89
$\delta$	−20°39′43″.0	+03°17′38″.7
$m_B$	17.46 ± 0.30 <sup>a</sup>	13.14 ± 0.30 <sup>a</sup>
$m_V$	12.19 ± 0.30 <sup>b</sup>	12.29 ± 0.10 <sup>b</sup>
$m_R$	11.85 ± 0.30 <sup>a</sup>	12.67 ± 0.30 <sup>a</sup>
$m_{r'}$	12.29 ± 0.30 <sup>c</sup>	12.36 ± 0.30 <sup>c</sup>
$m_I$	11.653 ± 0.020 <sup>d</sup>	11.326 ± 0.053 <sup>b</sup>
$m_J$	10.959 ± 0.026 <sup>e</sup>	11.263 ± 0.027 <sup>e</sup>
$m_H$	10.474 ± 0.022 <sup>e</sup>	10.913 ± 0.024 <sup>e</sup>
$m_K$	10.390 ± 0.023 <sup>e</sup>	10.892 ± 0.026 <sup>e</sup>
$T_{\text{eff}}$ (K)	4950 ± 60	5890 ± 120
$\log g_*$ (km s <sup>−1</sup> )	4.40 ± 0.20	4.27 ± 0.09
$\xi_t$ (km s <sup>−1</sup> )	0.60 ± 0.25	0.94 ± 0.15
$v \sin i_*$ (km s <sup>−1</sup> )	<2.7 ± 0.3	1.20 ± 0.73
[Fe/H]	0.22 ± 0.11	−0.36 ± 0.14
[Ca/H]	0.16 ± 0.15	−0.25 ± 0.09
[Sc/H]	0.19 ± 0.11	−0.18 ± 0.18
[Ti/H]	0.26 ± 0.15	−0.14 ± 0.11
[V/H]	0.44 ± 0.20	−0.29 ± 0.12
[Cr/H]	0.22 ± 0.11	−0.40 ± 0.11
[Mn/H]	0.29 ± 0.20	−0.55 ± 0.09
[Co/H]	0.24 ± 0.11	−0.36 ± 0.14
[Ni/H]	0.20 ± 0.12	−0.34 ± 0.15
$\log A(\text{Li})$	<0.32 ± 0.16	1.21 ± 0.10
Mass ( $M_\odot$ )	0.87 ± 0.08	1.04 ± 0.09
Radius ( $R_\odot$ )	0.96 ± 0.24	1.24 ± 0.15
Spectral type	K3	G1
Distance (pc)	235 ± 55	410 ± 70

*Note.* Mass and radius estimate using the Torres, Andersen & Giménez (2010) calibration. Spectral type estimated from  $T_{\text{eff}}$  using the table in Gray (2008). Abundances are relative to the solar values obtained by Asplund et al. (2009).

<sup>a</sup>NOMAD (Zacharias et al. 2004); <sup>b</sup>TASS (Droege et al. 2006); <sup>c</sup>CMC14, ViZier I/304/out; <sup>d</sup>DENIS (DENIS Consortium 2005); <sup>e</sup>2MASS (Skrutskie et al. 2006).

For future reference, we provide here the position of the visual companion. *Gaia* will soon produce parallaxes and proper motions that should confirm our analysis. If instead they are found at the same distance, implying they are gravitationally bound, then the companion must be an M dwarf with a much redder  $B - V$  than the one we measured.

### 3.2 TRAPPIST

Both WASP systems were observed with the 0.6-m Transiting Planets and Planetesimals Small Telescope (TRAPPIST) robotic telescope, also located at La Silla. Three transits were obtained on WASP-53 and seven on WASP-81. TRAPPIST is equipped with a thermoelectrically cooled 2K × 2K CCD, which has a pixel scale of 0.6 arcsec that translates into a 22 × 22 arcmin<sup>2</sup> field of view. For details of TRAPPIST, see Gillon et al. (2011) and Jehin et al. (2011). Two filters were used: a blue-blocking filter that has a transmittance of >90 per cent from 500 to beyond 1000 nm and an ‘I + z’ filter that has a transmittance of >90 per cent from 750 to beyond 1100 nm. During the runs, the positions of the stars on the chip were maintained to within a few pixels thanks to a ‘software guiding’ system that regularly derives an astrometric solution for the most recently acquired image and sends pointing corrections to the mount if needed. After a standard pre-reduction (bias, dark and flat-field correction), the stellar fluxes were extracted from

the images using the IRAF/DAOPHOT<sup>2</sup> aperture photometry software (Stetson 1987). For each light curve we tested several sets of reduction parameters and kept the one giving the most precise photometry for the stars of similar brightness as the target. After a careful selection of reference stars, the transit light curves were finally obtained using differential photometry. Some light curves were affected by a meridian flip, i.e. the 180° rotation that the German equatorial mount telescope has to undergo when the meridian is reached. This movement results in different positions of the stellar images on the detector before and after the flip, and thus in a possible jump of the differential photometry at the time of the flip. We have accounted for this in our light-curve analysis by including a normalization offset in our model at the time of the flip (see Figs C1 and C2 and Table C1). More details on data acquisition and data reduction can be found in Gillon et al. (2013) and Delrez et al. (2014), whose procedures were followed here as well. The times of observations are provided in JD(UTC), changed later to BJD(TDB) during the global analysis.

### 3.3 New Technology Telescope

One transit of WASP-53 was observed using the ESO Faint Object Spectrograph and Camera, v.2 (EFOSC2) instrument on the New Technology Telescope (NTT) at ESO’s observation of La Silla (ProgID 088.C-0204, PI: Tregloan-Reed; see Tregloan-Reed & Southworth 2013 for further details of this observing run). The instrument has a 2K × 2K CCD covering a 4.1 × 4.1 arcmin<sup>2</sup> field of view and a pixel scale of 0.12 arcsec. Observations were curtailed shortly after the end of the transit due to the onset of day-time. We observed through a Gunn-*r* filter (ESO filter #784) with heavy defocussing and exposure times of 150 s.

The data were reduced using the DEFOT pipeline (Southworth et al. 2009), which utilizes an aperture-photometry routine APER.PRO ported from DAOPHOT (Stetson 1987). The radius of the inner aperture was 45 pixels and the sky annulus extended from 60 to 100 pixels. Debiassing and flat-fielding the data did not make a significant difference to the results, so we neglected these calibrations.

A differential-photometry light curve was obtained for WASP-53 versus an ensemble comparison star. Due to the small field of view of NTT/EFOSC2, we were able to use only three comparison stars, all of which were at least two magnitudes fainter than WASP-53 in the *r* band. The weights of the comparison stars, used in summing their fluxes to create the ensemble comparison star, were optimized to minimize the scatter in the data outside transit. Finally, the time stamps were moved to the BJD(TDB) time-scale using routines from Eastman, Siverd & Gaudi (2010).

## 4 SPECTROSCOPIC OBSERVATIONS

We collected 98 CORALIE spectra on WASP-53 between 2010-12-04 and 2016-11-21, as well as 83 HARPS spectra between 2011-08-28 and 2014-09-28. On WASP-81, we acquired 67 spectra with CORALIE from 2011-09-28 to 2015-07-08, and 32 spectra using HARPS between 2013-04-20 and 2016-10-21.

<sup>2</sup> IRAF is distributed by the National Optical Astronomy Observatory, which is operated by the Association of Universities for Research in Astronomy, Inc., under cooperative agreement with the National Science Foundation.

#### 4.1 Spectral analysis

The analysis was performed on the standard pipeline reduction products using the methods given in Doyle et al. (2013). The  $H\alpha$  line was used to give an initial estimate of the effective temperature ( $T_{\text{eff}}$ ). The surface gravity ( $\log g_*$ ) was determined from the  $\text{Ca I}$  line at 6439 along with the  $\text{Na I D}$  lines. Additional  $T_{\text{eff}}$  and  $\log g_*$  diagnostics were performed using the Fe lines. An ionization balance between Fe I and Fe II was required along with a null dependence of the abundance on either equivalent width or excitation potential. This null dependence was also required to determine the microturbulence ( $\xi_t$ ). The parameters obtained from the analysis are listed in Table 2. Some of those parameters will later be employed as priors to compute the most likely physical parameters of each system. The elemental abundances were determined from equivalent width measurements of several clean and unblended lines, and additional least-squares fitting of lines was performed when required. The quoted error estimates include that given by the uncertainties in  $T_{\text{eff}}$ ,  $\log g_*$  and  $\xi_t$ , as well as the scatter due to measurement and atomic data uncertainties.

The individual HARPS spectra of WASP-53 were co-added to produce a single spectrum with an average signal-to-noise ratio (S/N) in excess of 100:1. The macroturbulence was assumed to be zero since for mid-K stars it is expected to be lower than that of thermal broadening (Gray 2008). The projected stellar rotation velocity ( $v \sin i_*$ ) was determined by fitting the profiles of several unblended lines, yielding an upper limit of  $2.7 \pm 0.3 \text{ km s}^{-1}$  for WASP-53. There is no significant detection of lithium in WASP-53. The equivalent width upper limit of  $13 \text{ mÅ}$  corresponds to an abundance upper limit of  $\log A(\text{Li}) < 0.32 \pm 0.16$ . This implies an age of at least several hundreds megayears (Sestito & Randich 2005).

Similarly, the combination of the HARPS spectra obtained on WASP-81 produced a combined spectrum with an average S/N of 60:1. Here, we used a macroturbulent value of  $3.84 \pm 0.73 \text{ km s}^{-1}$  after a relation from Doyle et al. (2014).  $v \sin i_*$  was found to be  $1.20 \pm 0.69 \text{ km s}^{-1}$ .

#### 4.2 Radial velocities

The spectra were reduced using the standard CORALIE and HARPS reduction software. They have been shown to reach remarkable precision and accuracy, reaching below  $1 \text{ m s}^{-1}$  (e.g. Lovis et al. 2006; Marmier et al. 2013; López-Morales et al. 2014). We extracted the radial velocities for both stars by cross-correlating each spectrum with a binary mask. For WASP-53A we used a mask corresponding to a K5 spectral type, and for WASP-81A we employed a G2 mask. We fit the corresponding cross-correlation function with a Gaussian, whose mean provides us with the radial velocity (Baranne et al. 1996). The corresponding values are displayed in the Journal of Observations (Appendix A) along with observational data such as the individual exposure times. The observations are graphically presented in Fig. 5a for WASP-53 and in Fig. 5b for WASP-81.

##### 4.2.1 WASP-53

We first started to monitor WASP-53A with CORALIE on 2010 December 5 targeting the orbital phase 0.75. The second spectrum was obtained 28 nights later, close to phase 0.25, and revealed a blueshifted movement of nearly  $400 \text{ m s}^{-1}$ , compatible with a planetary object. The star was immediately flagged for intense follow-up, with the third spectrum acquired two nights after the second. The star was moving rapidly, and there was no sign of a change in

the line width [full width at half-maximum (FWHM)], nor of its shape (span of the bisector slope), as can be visually inspected in Figs 5 and 6. However, despite being obtained at phase 0.81, the star's velocity was puzzling, being redshifted by  $300 \text{ m s}^{-1}$  compared to the first epoch. Our strategy has always been to follow any radial velocity movement and identify its origin, planetary or not. Observations were continued.

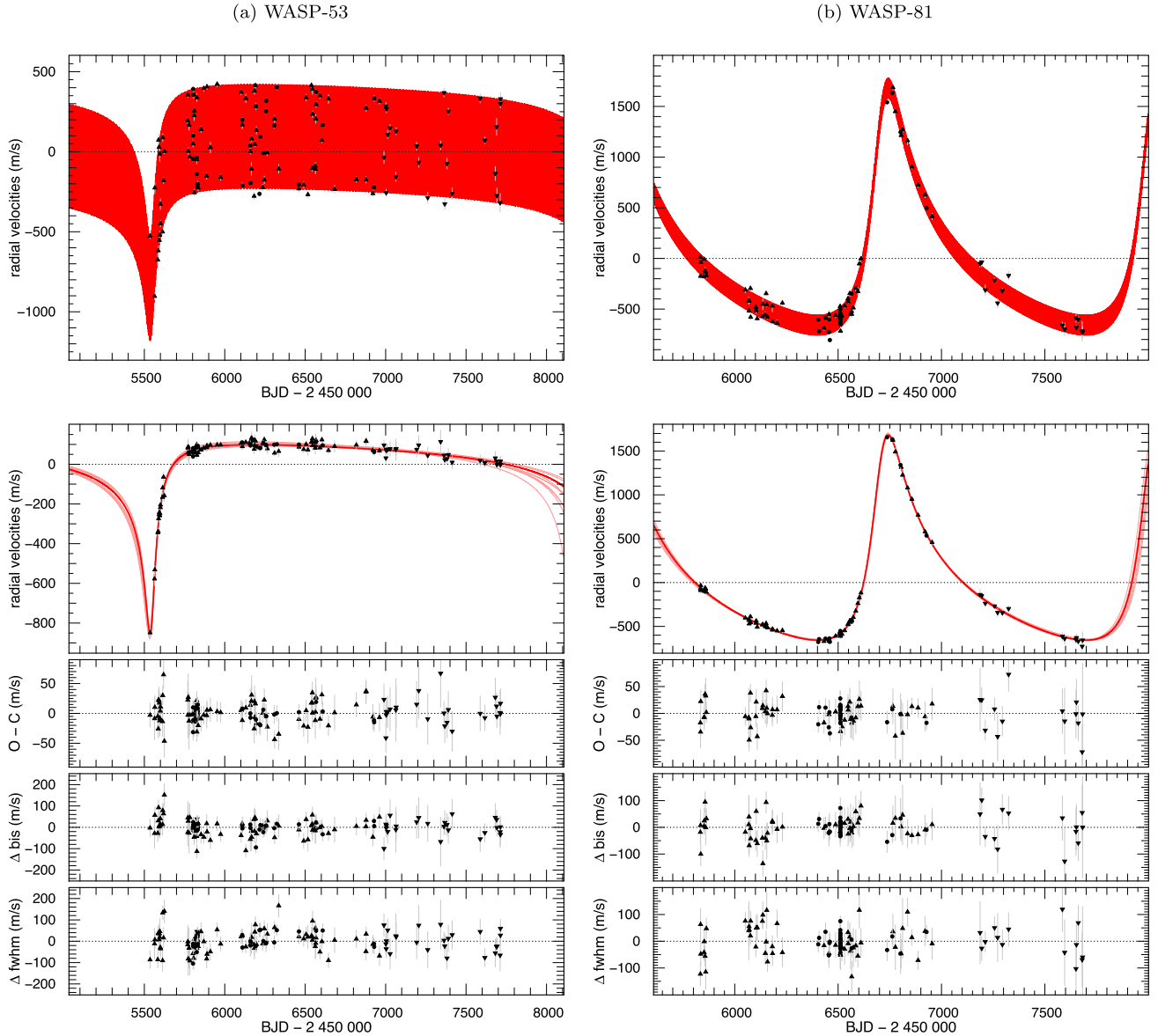
Before the observing season was over, we had confirmed a radial velocity oscillation at a period matching the WASP signal (caused by WASP-53b) plus a rapid rise in the radial velocity. The following season we observed with both HARPS and CORALIE. WASP-53b's motion was quickly recovered and found to be in phase. Our monitoring continued and we observed the velocity of the star rise until it plateaued, indicating that an additional, massive, highly eccentric object had just finished passing through periastron. Our observations are ongoing and, to this day, the velocity of the star has yet to decrease to the level observed at our first spectrum. In total we present 98 spectra with CORALIE, including 25 since an upgrade that saw the installation of new, octagonal, fibres (2014 November) and of a Fabry–Pérot (2015 April) for the wavelength calibration throughout the night. We also gathered 83 spectra with HARPS, which include three time series obtained during transit, in order to capture the Rossiter–McLaughlin effect (Fig. 7).

The CORALIE data were divided into two independent data sets in order to account for any offset between before and after the upgrade. The HARPS set was divided into four sets, one for each Rossiter–McLaughlin effect (plus the measurement obtained the night before and after transit) and one set containing all the rest of the data. We obtained similar results by analysing the HARPS data as one set; however, the division mitigates against any activity effect, such as done in Triaud et al. (2009). The Journal of Observations (Appendix A) is separated in several tables according to the various subsamples of radial velocities.

##### 4.2.2 WASP-81

CORALIE acquired our first spectrum on the WASP-81 system on 2011 September 29. The following night another was observed at the opposite phase. Its radial velocity revealed a variation and the target was placed in high priority. Within a month we had confirmed a variation at the WASP photometric period, and the star had set. The following season, we intended to monitor just enough to confirm that the oscillation was still in phase and to start routine long-term monitoring, as in the case of WASP-47 (Neveu-VanMalle et al. 2016). The first measurement had a value clearly below expectations so we resumed an intense follow-up. As with WASP-53, we requested observing time on HARPS and monitored the system in parallel with CORALIE. On the third season, as we had predicted, the velocity reached a minimum and started rising. HARPS was the first instrument on sky for the fourth season. We had anticipated that the system would have returned to a similar velocity as in the first points. We were surprised to find that the star was nearly  $2 \text{ km s}^{-1}$  higher than in the previous season. Shortly after, velocities started to drop and the outer orbit closed earlier this year. In total 67 spectra were collected with CORALIE, 14 of which were after the upgrade. With HARPS we gathered 32 measurements, of which 19 were obtained during a single night as an attempt to detect the Rossiter–McLaughlin effect (Fig. 7).

As with WASP-53 we separate the CORALIE data into two sets, before and after upgrade. Since we do not detect the Rossiter–McLaughlin effect, all of the HARPS points are analysed as part of



**Figure 5.** Radial velocities and models for WASP-53 (left) and WASP-81 (right). HARPS points are represented by discs, CORALIE is shown as triangles, and pointed upwards prior to the upgrade and downwards for data acquired since. Top: radial velocity time series with the preferred two-planet model adjusted to the data. Middle: same as on top, minus the inner planet. Ten alternative models sampled randomly from the posterior are displayed in pink. Rossiter–McLaughlin sequences were removed from these plots; they can be found in Fig. 7. Residuals are shown below the main plots. Further down, in order, we have the variation in the slope of the bisector span, and the variation in the FWHM of the cross-correlation function.

the same set. Refer to the Journal of Observations (Appendix A) for further details.

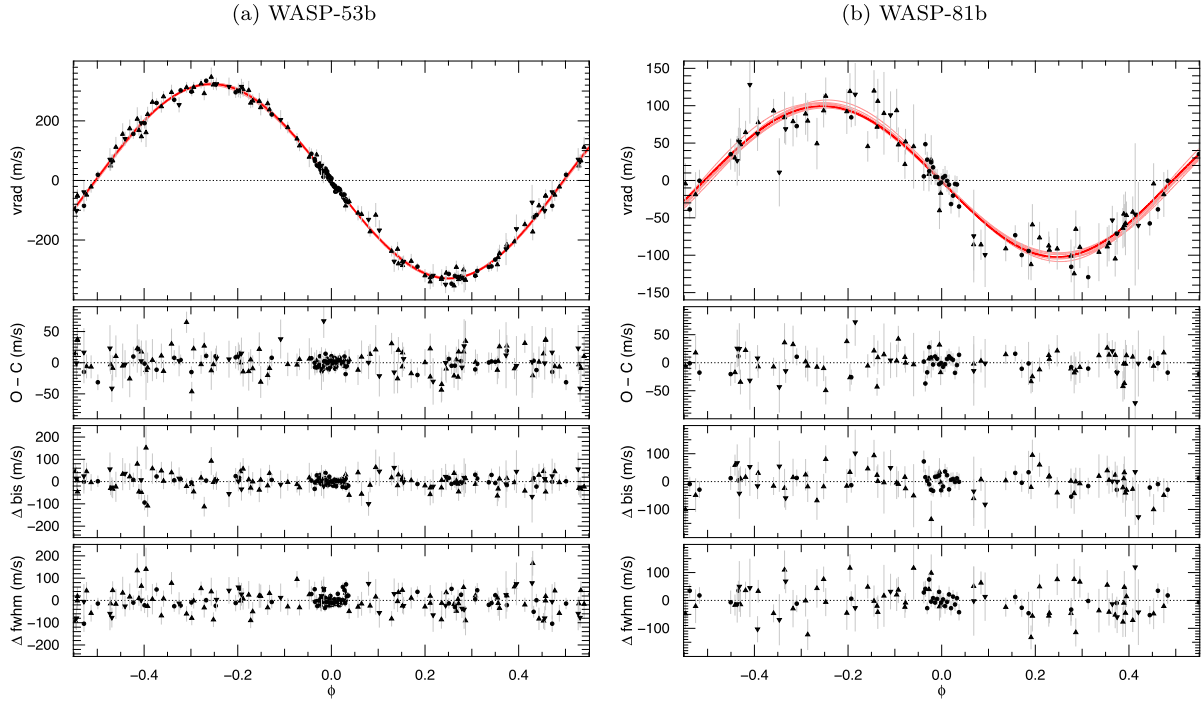
## 5 GLOBAL ANALYSIS

We analysed all of the photometric data and all of the radial velocity data together. We estimated the observed and physical parameters of the system using a Markov chain Monte Carlo (MCMC) algorithm, as detailed in Gillon et al. (2012). Our modus operandi is similar with the difference described below.

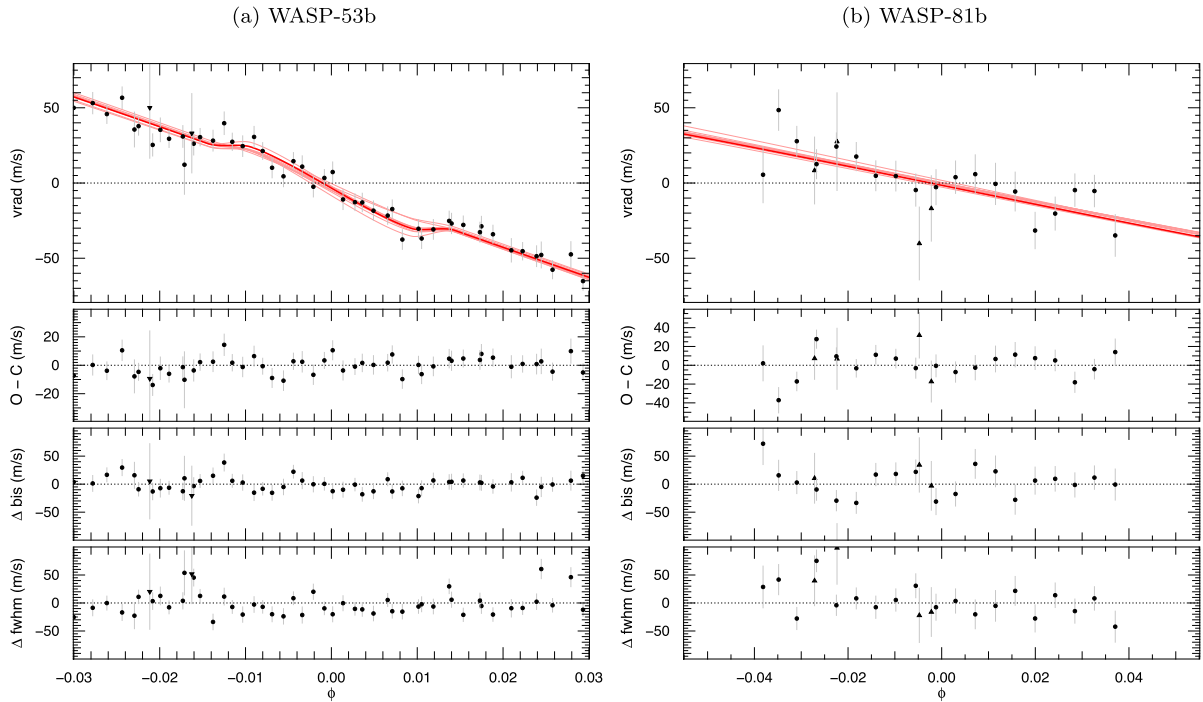
The photometric light curves were modelled using the formalism of Mandel & Agol (2002) for a transiting planet. The TRAPPIST and EulerCam light curve time stamps were transformed into  $\text{BJD}_{\text{TDB}}$  from  $\text{JD}_{\text{UTC}}$ . The limb darkening was included in the

form of a quadratic law, and its parameters were allowed to float, within the constraints of priors. The priors were computed by interpolating the data tabulated by Claret (2004), consistent with the stellar parameters of Table 2. In addition, on each photometric time series, we allowed for a quadratic polynomial as a function of time in order to adjust for differential extinction relative to the ensemble of comparison stars. After this treatment, some light curves still contained a significant residual scatter. For these we tried a number of other detrending functions, selecting them on the basis of a reduction in the global Bayesian information criterion (BIC; Schwarz 1978). Those comparisons were performed by systematically using the same starting seed for the random elements. Details about which functions were selected and applied can be found in Appendix C. For those light curves where any new degree of





**Figure 6.** Radial velocity measurements phased with the orbit of the inner planet, after subtracting the variation due to the outer object. The preferred model is drawn in thick red, with another 10 alternate models randomly picked from the posterior shown in pink. Residuals are shown below the main plots. Further down we have the variation in the slope of the bisector span and the variation in the FWHM of the cross-correlation function. The symbols are the same as in Fig. 5.



**Figure 7.** Same as for Fig. 6, but zoomed around transit time to show the Rossiter–McLaughlin effect.

complexity led to a worsened BIC, but where nevertheless the overall  $\chi^2$  implied a poor fit, we scaled<sup>3</sup> our error bars so as

<sup>3</sup> We realize that this approach is not self-consistent, but our method is sufficient for small re-adjustments of the uncertainties and unambiguous detections.

to approach a general reduced  $\chi_r^2 = 1$ . The uncertainty increase ensures that our confidence intervals are not underestimated. Those correction factors are available in Table C1. A similar approach is used for the radial velocities, by quadratically adding a *jitter* term to some sequences. Two of the nine sequences required a jitter of order  $1.5\text{--}2\text{ m s}^{-1}$ .



Individual detrended light curves and their transit model can be visually inspected in Figs 3 and 4. Residuals and fully modelled light curves are presented in Appendix C.

The radial velocities around the orbit were adjusted with two eccentric Keplerian functions (as in Hilditch 2001), neglecting Newtonian effects. Since we cover less than an orbital period for WASP-53c and barely one for WASP-81c, we did not include a long-term trend. Its inclusion leads to higher BIC values and its slope is mostly unconstrained. The Rossiter–McLaughlin effect was computed using the code written by Giménez (2006), following the formalism of Kopal (1942). For WASP-81b, the effect was not detected and the final fit does not include it. For both systems, we find the inner planets’ orbits to be consistent with circular. Although allowing for eccentricity adds two parameters and increases the BIC, we nevertheless let these parameters float so as to include their uncertainties when marginalizing the other parameters.

The MCMC’s jump parameters are mostly set to match observables, which are then converted to physical parameters to compute the relevant models.  $D$  is the transit depth,  $b$  the impact parameter,  $W$  the transit width,  $T_0$  the mid-transit time and  $P$  is the orbital period. We combine the eccentricity  $e$  and the angle of periastron  $\omega$  into the pair  $\sqrt{e} \cos \omega$ ,  $\sqrt{e} \sin \omega$  which helps when exploring small eccentricities (Triaud et al. 2011). Similarly we also construct the pair  $\sqrt{v \sin I_*} \cos \beta$ ,  $\sqrt{v \sin I_*} \sin \beta$  to model the Rossiter–McLaughlin effect.  $\beta$  is the projected spin–orbit angle, and  $v \sin I_*$  is the measure of the projected rotation velocity of the star, which in principle should match  $v \sin i_*$  from the spectral analysis. Instead of the semi-amplitude  $K$  we use the jump parameter  $K_2$  such that  $K_2 = K \sqrt{1 - e^2} P^{1/3}$ . This helps reduce some correlation between the parameters (Ford 2006) helping with the exploration of parameter space. For similar reasons, we combine the limb-darkening coefficients into  $c_1 = 2u_1 + u_2$  and  $c_2 = u_1 - 2u_2$  following the recommendation of Holman et al. (2006). In the case of two *Euler* light curves for WASP-53 we also fit a sine function through the data, with a period  $P$  and a  $T_0$ . Additional subscript indicates which light curves those are for.

All jump parameters are well constrained within one scale. Except for  $P_c$  and  $K_{2,c}$ , we sample our parameters’ posteriors using Gaussian priors, whose variance is set with an initial Gibbs sampler. As an improvement over Gillon et al. (2012)  $P_c$  and  $K_{2,c}$  were sampled using non-informative priors in log space, otherwise known as *Jeffrey* priors. However their values remained largely within one error bar, and the application of Gaussian steps does not lead to qualitatively different results. The metallicity [Fe/H] and effective temperature  $T_{\text{eff}}$  complete this list of jump parameters. They are controlled by priors obtained from Table 2 and used to compute at every MCMC step a stellar mass and a stellar radius produced in a fashion similar to Torres et al. (2010).

For our final analysis we set 10 chains of 100 000 steps, starting from different seeds. All converged to similar BIC values. The first 20 000 steps were systematically removed (to allow for burn-in), and the remainder were analysed leading to our results. For each of our 10 chains, we extract the median value for each parameter and compare them to one another. They are usually of order 0.1 per cent different from one another, except for the jump parameters responsible for modelling WASP-53c, which can vary as much as 60 per cent. This is because the orbit is not closed. We discuss this further in the next section.

The posterior probability distributions have been stored and can be requested by e-mail to the lead author. We present the median values and  $1\sigma$  region of our posteriors in Table 3 for the jump

**Table 3.** Median and  $1\sigma$  confidence regions for the jump parameters that evolve in our MCMC chains. Errors on the last two digits of each parameter are given in brackets. Parameters not dependent on information contained within the adjusted data are highlighted with asterisks.

Parameters (units)	WASP-53	WASP-81	
The star			
[Fe/H] (dex)	0.22 <sup>(+0.11)</sup> <sub>(−0.11)</sub>	−0.36 <sup>(+0.14)</sup> <sub>(−0.14)</sub>	*
$T_{\text{eff}}$ (K)	4953 <sup>(+59)</sup> <sub>(−60)</sub>	5870 <sup>(+120)</sup> <sub>(−120)</sub>	*
$c_{1,IR}$	—	0.918 <sup>(+33)</sup> <sub>(−33)</sub>	*
$c_{2,IR}$	—	−0.280 <sup>(+23)</sup> <sub>(−23)</sub>	*
$c_{1,BB}$	1.069 <sup>(+37)</sup> <sub>(−38)</sub>	0.772 <sup>(+57)</sup> <sub>(−57)</sub>	*
$c_{2,BB}$	0.031 <sup>(+31)</sup> <sub>(−30)</sub>	−0.321 <sup>(+37)</sup> <sub>(−38)</sub>	*
$c_{1,I+z}$	0.995 <sup>(+62)</sup> <sub>(−61)</sub>	0.807 <sup>(+55)</sup> <sub>(−57)</sub>	*
$c_{2,I+z}$	−0.020 <sup>(+37)</sup> <sub>(−37)</sub>	−0.297 <sup>(+30)</sup> <sub>(−32)</sub>	*
$c_{1,r'}$	1.316 <sup>(+29)</sup> <sub>(−29)</sub>	—	*
$c_{2,r'}$	0.272 <sup>(+28)</sup> <sub>(−28)</sub>	—	*
Inner planet			
$D_b$	0.01831 <sup>(+33)</sup> <sub>(−34)</sub>	0.01254 <sup>(+27)</sup> <sub>(−26)</sub>	
$b_b$ ( $R_\odot$ )	0.562 <sup>(+20)</sup> <sub>(−22)</sub>	0.15 <sup>(+0.10)</sup> <sub>(−0.10)</sub>	
$W_b$ (d)	0.09469 <sup>(+65)</sup> <sub>(−64)</sub>	0.14501 <sup>(+73)</sup> <sub>(−64)</sub>	
$T_{0,b}$ (BJD − 245 0000)	5943.56695 <sup>(+11)</sup> <sub>(−12)</sub>	6195.57462 <sup>(+20)</sup> <sub>(−20)</sub>	
$P_b$ (d)	3.3098443 <sup>(+20)</sup> <sub>(−20)</sub>	2.7164762 <sup>(+23)</sup> <sub>(−23)</sub>	
$\sqrt{e_b} \cos \omega_b$	−0.070 <sup>(+24)</sup> <sub>(−21)</sub>	−0.026 <sup>(+83)</sup> <sub>(−76)</sub>	
$\sqrt{e_b} \sin \omega_b$	−0.085 <sup>(+46)</sup> <sub>(−33)</sub>	−0.00 <sup>(+0.12)</sup> <sub>(−0.12)</sub>	
$K_{2,b}$ ( $\text{m s}^{-1} \text{d}^{1/3}$ )	485.9 <sup>(+2.7)</sup> <sub>(−2.7)</sub>	140.7 <sup>(+4.7)</sup> <sub>(−4.6)</sub>	
$\sqrt{v \sin I_*} \cos \beta_b$	0.91 <sup>(+0.11)</sup> <sub>(−0.13)</sub>	—	
$\sqrt{v \sin I_*} \sin \beta_b$	−0.07 <sup>(+0.19)</sup> <sub>(−0.19)</sub>	—	
Outer planet			
$T_{0,c}$ (BJD − 245 0000)	5456 <sup>(+11)</sup> <sub>(−13)</sub>	6936.5 <sup>(+2.6)</sup> <sub>(−2.5)</sub>	
$P_c$ (d)	2840 <sup>(+170)</sup> <sub>(−130)</sub>	1297.2 <sup>(+8.1)</sup> <sub>(−7.8)</sub>	
$\sqrt{e_c} \cos \omega_c$	−0.867 <sup>(+12)</sup> <sub>(−12)</sub>	0.5871 <sup>(+40)</sup> <sub>(−42)</sub>	
$\sqrt{e_c} \sin \omega_c$	−0.292 <sup>(+30)</sup> <sub>(−30)</sub>	−0.4609 <sup>(+61)</sup> <sub>(−61)</sub>	
$K_{2,c}$ ( $\text{m s}^{-1} \text{d}^{1/3}$ )	3685 <sup>(+110)</sup> <sub>(−89)</sub>	10 557 <sup>(+51)</sup> <sub>(−50)</sub>	
Jump parameters	23	21	

parameters and in Table 4 for the physical parameters. The results are discussed in the following section.

## 6 RESULTS

### 6.1 WASP-53

WASP-53 is a system composed of a star and two orbiting objects. WASP-53b is a hot Jupiter with a mass  $M_b = 2.1 \pm 0.1 M_{\text{Jup}}$  and a radius  $R_b = 1.07 \pm 0.04 R_{\text{Jup}}$ . We find the inner orbit to be consistent with zero eccentricity, placing a 99 per cent confidence limit of 0.03. The Rossiter–McLaughlin effect is weakly detected. We find a lower amplitude than we anticipated, likely because  $v \sin i_*$  is overestimated owing to an underestimation of macroturbulence, as already noted for a number of late-type dwarfs (e.g. Triaud et al. 2011, 2015). We find the spin–orbit angle  $\beta = -4^\circ \pm 12^\circ$ . The planet appears to be coplanar.

The orbit of WASP-53c does not close within the span of our observations, which affects the various chains we launched. The

**Table 4.** Final estimates for the median and  $1\sigma$  confidence regions, for various interesting parameters of the WASP-53 and WASP-81 systems. They were estimated from the posteriors of the jump parameters outlined in Table 3. Errors on the last two digits of each parameter are given in brackets. Upper limits are for  $3\sigma$  confidence. Parameters not dependent on information contained within the adjusted data are highlighted with asterisks.

Physical parameters	WASP-53	WASP-81	
The star			
$M_*$ ( $M_\odot$ )	$0.839^{(+54)}_{(-54)}$	$1.080^{(+59)}_{(-58)}$	*
$R_*$ ( $R_\odot$ )	$0.798^{(+23)}_{(-23)}$	$1.283^{(+40)}_{(-37)}$	
$\rho_*$ ( $\rho_\odot$ )	$1.648^{(+91)}_{(-85)}$	$0.513^{(+30)}_{(-37)}$	
$L_*$ ( $L_\odot$ )	$0.344^{(+23)}_{(-23)}$	$1.76^{(+0.20)}_{(-0.18)}$	
$T_{\text{eff}}$ (K)	$4953^{(+60)}_{(-60)}$	$5870^{(+120)}_{(-120)}$	*
$\log g_*$ (cgs)	$4.553^{(+19)}_{(-20)}$	$4.258^{(+22)}_{(-27)}$	*
[Fe/H] (dex)	$0.22^{(+0.11)}_{(-0.11)}$	$-0.36^{(+0.14)}_{(-0.14)}$	*
$v \sin I_*$ ( $\text{km s}^{-1}$ )	$0.86^{(+0.21)}_{(-0.21)}$	—	
Inner planet			
$P_b$ (d)	$3.3098443^{(+20)}_{(-20)}$	$2.7164762^{(+23)}_{(-23)}$	
$T_{0,b}$ (BJD – 245 0000)	$5943.56695^{(+11)}_{(-12)}$	$6195.57462^{(+23)}_{(-20)}$	
$K_b$ ( $\text{m s}^{-1}$ )	$326.1^{(+1.8)}_{(-1.8)}$	$100.8^{(+3.4)}_{(-3.3)}$	
$M_b$ ( $M_{\text{Jup}}$ )	$2.132^{(+92)}_{(-94)}$	$0.729^{(+36)}_{(-35)}$	
$R_b$ ( $R_{\text{Jup}}$ )	$1.074^{(+37)}_{(-37)}$	$1.429^{(+51)}_{(-46)}$	
$\rho_b$ ( $\rho_{\text{Jup}}$ )	$1.72^{(+0.15)}_{(-0.13)}$	$0.250^{(+23)}_{(-23)}$	
$\log g_b$ (cgs)	$3.680^{(+22)}_{(-22)}$	$2.967^{(+27)}_{(-30)}$	
$a_b/R_*$	$11.05^{(+0.20)}_{(-0.19)}$	$6.56^{(+0.13)}_{(-0.16)}$	
$T_{b,\text{eq}}$ (K)	$1053^{(+16)}_{(-16)}$	$1623^{(+38)}_{(-37)}$	
$a_b$ (au)	$0.04101^{(+83)}_{(-91)}$	$0.03908^{(+70)}_{(-72)}$	
$i_b$ ( $^\circ$ )	$87.08^{(+0.16)}_{(-0.15)}$	$88.69^{(+0.88)}_{(-0.92)}$	
$\beta_b$ ( $^\circ$ )	$-4^{(+12)}_{(-12)}$	—	
$e_b$	$<0.030$	$<0.066$	
Outer planet			
$P_c$ (d)	$>2840^{(+170)}_{(-130)}$	$1297.2^{(+8.1)}_{(-7.8)}$	
$T_{0,c}$ (BJD – 245 0000)	$5456^{(+11)}_{(-13)}$	$6936.5^{(+2.6)}_{(-2.5)}$	
$K_c$ ( $\text{m s}^{-1}$ )	$>475.6^{(+8.2)}_{(-8.0)}$	$1169.3^{(+6.9)}_{(-6.6)}$	
$M_c \sin i_c$ ( $M_{\text{Jup}}$ )	$>16.35^{(+0.85)}_{(-0.82)}$	$56.6^{(+2.0)}_{(-2.0)}$	
$a_c$ (au)	$>3.73^{(+0.16)}_{(-0.14)}$	$2.426^{(+44)}_{(-45)}$	
$e_c$	$0.8369^{(+69)}_{(-70)}$	$0.5570^{(+44)}_{(-44)}$	
$\omega_c$ ( $^\circ$ )	$198.6^{(+2.0)}_{(-2.0)}$	$321.86^{(+0.52)}_{(-0.51)}$	

median values on individual jump parameters values can vary by as much as 50 per cent, which is reflected in the large errors in Tables 3 and 4. This means that while we provide median values and their  $1\sigma$  confidence ranges, those are in fact more akin to lower limits on  $M_c$ ,  $e_c$ ,  $P_c$ , etc. WASP-53c is at least  $16M_{\text{Jup}}$ , with a period likely longer than 2500 d. The eccentricity of its orbit is high with our data being most consistent with  $0.84 \pm 0.01$ .

We were lucky to observe the system during the final phases of WASP-53c’s periastron passage (but unlucky to miss the first half). If removing the first season of CORALIE data, we only detect a quadratic drift with a weak curvature and would never have guessed the presence of such a massive companion within the system. We searched for transits of WASP-53c and WASP-81c around the dates

**Table 5.** Dates on which WASP-53c and WASP-81c may transit. Numbers are calendar dates and Barycentric Julian Dates (BJD) – 245 0000.

WASP-53c	WASP-81c
Passed dates	
—	$3044.9^{+24}_{-24}$ (2004 February 9)
—	$4342.1^{+16}_{-16}$ (2007 August 29)
$2616.0^{+170}_{-130}$ (2002-12-07)	$5639.3^{+8.5}_{-8.2}$ (2011 March 18)
$5455.5^{+11}_{-13}$ (2010-09-16)	$6936.5^{+2.6}_{-2.5}$ (2014 October 6)
Future dates	
$8295.1^{+170}_{-130}$ (2018-06-25)	$8233.8^{+8.5}_{-8.2}$ (2018 April 25)
—	$9531.0^{+16}_{-16}$ (2021 November 12)

shown in Table 5, but could not see anything convincing. Future epochs are provided.

## 6.2 WASP-81

WASP-81 is a system composed of a star and two orbiting objects. WASP-81b is a hot Jupiter whose mass is  $M_b = 0.73 \pm 0.04M_{\text{Jup}}$  and radius is  $R_b = 1.43 \pm 0.05R_{\text{Jup}}$ . Its orbit is consistent with being circular and we place a 99 per cent confidence upper limit at 0.07. The Rossiter–McLaughlin effect is not detected. Its amplitude is projected to be less than  $10 \text{ m s}^{-1}$ . The low impact parameter means that the spin–orbit angle will be degenerate with  $v \sin I_*$  as in Triaud et al. (2011).

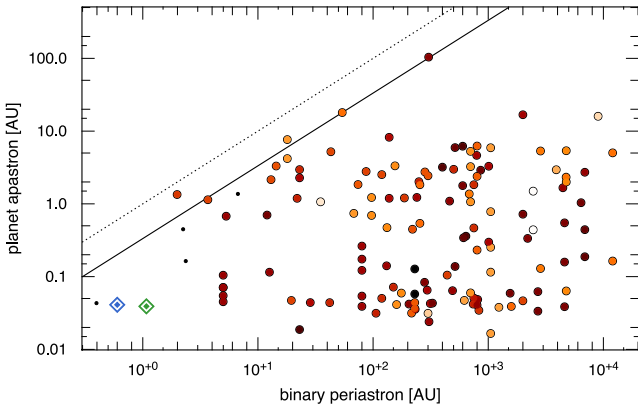
WASP-81c has a minimum mass  $M_c = 57 \pm 2M_{\text{Jup}}$ , a period  $P_c = 1297 \pm 8 \text{ d}$  and an eccentricity of order 0.56.

## 7 DISCUSSION

We have discovered two hot Jupiters orbiting the primary star of two tight binary systems. WASP-53b is super-Jupiter in mass, while WASP-81b is sub-Jupiter. Both occupy orbits that are typical for hot Jupiters (e.g. Santerne et al. 2016). Those two planets are both accompanied by brown dwarf mass objects, on highly eccentric orbits of a few au. In Fig. 8, we plot other known planetary systems orbiting one star of a multiple stellar system, showing how atypical WASP-53 and WASP-81 are within the current exoplanet population. Only four other systems have a stellar companion with periastra closer than 10 au: Kepler-444 (Dupuy et al. 2016), KOI-1257 (Santerne et al. 2014), HD 59686 (Ortiz et al. 2016) and the astonishing, maybe retrograde,  $\nu$  Oct (Ramm 2015) (the leftmost red dot, above the plain line). In addition, four gas giants have outer brown dwarf companions: HAT-P-13 (Knutson et al. 2014), HIP 5158 (Feroz, Balan & Hobson 2011), HD 168443 (Sahlmann et al. 2011) and HD 38529 (Benedict et al. 2010), in architectures similar to WASP-53 and WASP-81.

We now review the elements that make those systems stand out. We speculate about their origin and propose some observational tests to verify some of our scenarios.

We use the stability criterion numerically determined by Holman & Wiegert (1999) to compute the widest orbital separation that each of the hot Jupiters could have occupied. In the case of WASP-53 we obtain a critical semimajor axis  $a_{\text{crit}} = 0.16 \pm 0.15 \text{ au}$ , and for WASP-81,  $a_{\text{crit}} = 0.38 \pm 0.06 \text{ au}$ . This criterion was numerically determined for a mass ratio  $\mu = m_2/(m_1 + m_2) > 0.1$ , which is not satisfied for either of our systems and likely explains the large uncertainty for WASP-53. We therefore proceed with a stability



**Figure 8.** Planetary apastron versus binary periastron (or separation if eccentricity is unknown) in astronomical units for known S-type planetary systems. The colour of the dots reflects the logarithm of the ratio of the planet-hosting star mass to the mass of its stellar companion(s) (white = 0.2, black = 18). WASP-53 and WASP-81 are highlighted as a blue and green diamond, respectively. The small black dots represent four systems containing a gas giant and a brown dwarf. The dotted line is a 1:1 line, and the plain line is a 3:1 contour. Above that line systems are usually unstable (Dvorak 1986; Holman & Wiegert 1999). Data are collected from [openexoplanet-catalogue.com](http://openexoplanet-catalogue.com), <http://www.univie.ac.at/adg/schwarz/multiple.html> and [exoplanet.eu](http://exoplanet.eu).

criterion devised by Petrovich (2015b) for hierarchical planetary system. We find that unstable orbit start emerging for  $0.14 < a_b < 0.17$  for WASP-53, and  $0.30 < a_b < 0.38$  for WASP-81, in good agreement with the previous approach.

We further verify this with the method proposed by Cincotta, Giordano & Simó (2003), which uses a marker called the Mean Exponential Growth factor of Nearby Orbits (MEGNO), as implemented in *REBOUND* by Rein & Tamayo (2015). The MEGNO is a good tracker of orbital chaos, meaning that infinitesimal changes in initial parameters lead to diverging solutions. For quasi-periodic orbits, thus those showing no chaotic behaviour, the MEGNO reaches a value of 2 (Hinse et al. 2010). Larger values of the MEGNO, typically  $>4$ , indicate significant changes in the orbital parameters, a sign of chaos. Although this does not necessarily translate by unstable orbits (e.g. Deck et al. 2012), it often tracks them as we saw above. For our case, the MEGNO outlines where nearly closed orbits exist and therefore informs us on where any disc material may have been stable or where additional planets may exist.

We used the parameters provided in Table 4 for the outer companions, assumed coplanarity between the inner and outer orbits, and computed the MEGNO for a particle with a mass and separation for WASP-53b and WASP-81b. We integrated each system for  $5 \times 10^6$  yr and obtained values of 2.0062 and 1.9999, respectively, indicating stability. We expanded those simulation and explore the  $(a, e)$  parameter space. To compute the maps presented in Fig. 9 we integrated each of the pixels over  $5 \times 10^4$  yr. We observe from those results that wide regions of chaos exist, which is consistent with the work of Holman & Wiegert (1999) and Petrovich (2015b). Both WASP-53 and WASP-81 appear stable over long periods of time. In the case of WASP-53, we observe that only regions closer than 0.15 au retain stable orbits. For WASP-81, there is slightly more space with orbits within 0.3 au being generally stable.

We also investigated whether other planets could exist between objects b and c by adding a third massless particle. The results are displayed in Fig. 10, which show a range of stable orbits (in green). While this suggests that other planets could still be identified within

our systems, we think it unlikely since hot Jupiters are usually found to be isolated. Only one hot Jupiter is known to have other planetary companions within an astronomical unit (Becker et al. 2015; Neveu-VanMalle et al. 2016), and a recent analysis of the *Kepler* data shows that a lack of nearby companions is an important aspect that sets hot Jupiters apart from other gas giants (Huang, Wu & Triaud 2016).

## 7.1 Determination of the mass and orbital inclination of WASP-53c and WASP-81c

*Gaia* (Perryman et al. 2001) is an ESA mission launched in 2013 currently scanning the sky and measuring stellar positions with a precision of order  $30 \mu\text{as}$  for stars brighter than optical magnitude 12 (de Bruijne 2012). Several studies have investigated *Gaia*'s potential for detecting gas giants using astrometry (Casertano et al. 2008; Perryman et al. 2014; Sozzetti et al. 2014; Sahlmann, Triaud & Martin 2015), with Neveu et al. (2012) looking into the combination of radial velocities with astrometric measurements.

From Perryman et al. (2014), WASP-53c and WASP-81c will produce an astrometric displacement of their host stars,  $\alpha$ , defined as

$$\alpha = \left( \frac{M_p}{M_*} \right) \left( \frac{a}{1 \text{ au}} \right) \left( \frac{d}{1 \text{ pc}} \right)^{-1} \text{ arcsec.} \quad (1)$$

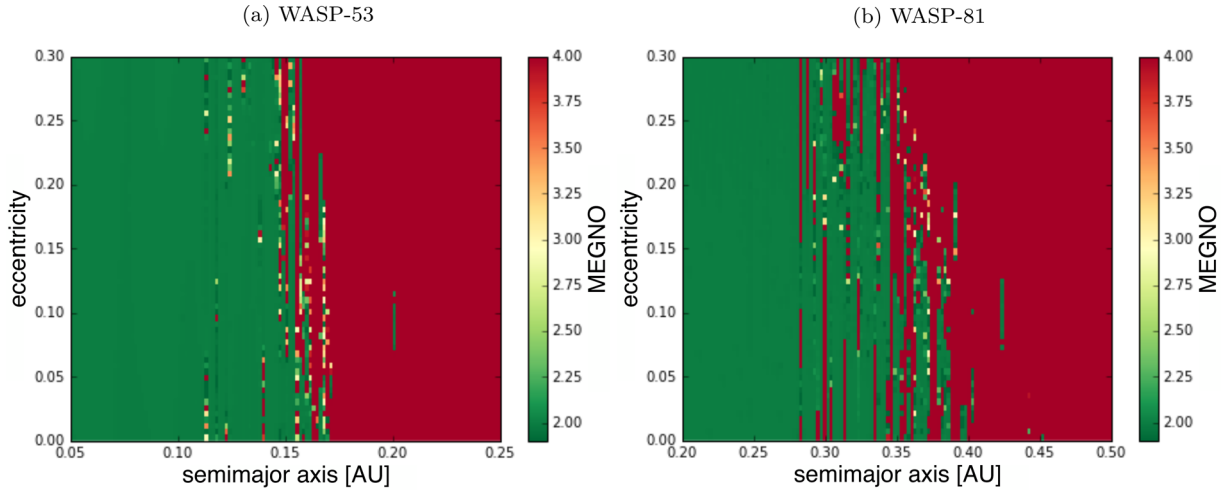
WASP-53 and WASP-81 will move on the sky by  $\alpha = 260$  and  $300 \mu\text{as}$ , respectively, caused by their outer companions, assuming an orbital inclination  $i_c = 90^\circ$ , which is the poorest scenario possible. We can expect of order  $N_{\text{obs}} = 70$  astrometric measurements<sup>4</sup> with typical uncertainties  $\sigma = 40 \mu\text{as}$  (de Bruijne 2012) to be collected on our two targets. This translates to an astrometric S/N, where  $S/N = \alpha \sqrt{N_{\text{obs}}} / \sigma$ , in excess of 50 for both systems. If instead  $i_c = 10^\circ$ , we obtain  $\alpha = 1500$  and  $1700 \mu\text{as}$ , respectively.<sup>5</sup> The amplitude of the orbital motion alone should inform us of the mutual inclination between the inner and the outer planet. For astrometric S/N values of  $>20$ , the orbital inclination will typically be estimated with a precision of  $<10^\circ$  (Sahlmann et al. 2015).

## 7.2 Transit-timing variations

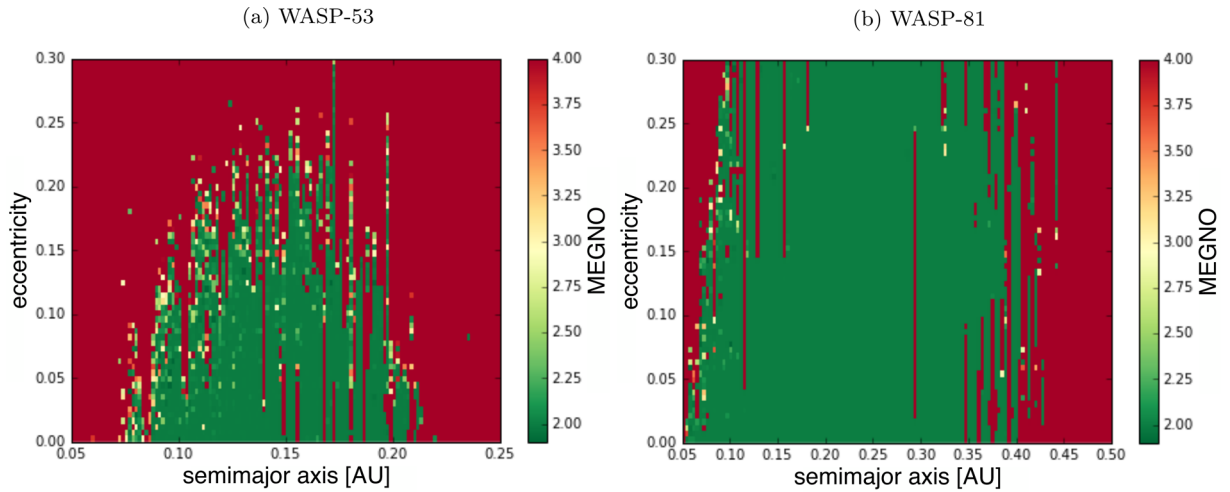
Using *REBOUND* (Rein & Spiegel 2015), we integrated the system over a few orbital periods of the outer companion and recorded when transits of the inner planet happened. For WASP-53b we expect total transit-timing variations, caused by the perturbing effect of the outer companion, to be of order 35 s. However, most of the variation happens during periastron, which means that during nearly 100 transit epochs we covered for WASP-53b, we expect no variation to be measurable. We expect a detectable signal to appear within the next 2 yr if our solution for the outer companion is correct. After WASP-53c swings again via periastron, the ephemeris for WASP-53b will become offset by approximately 35 s. This offset remains constant until the following periastron when it will offset again by the same amount. We repeated the procedure for WASP-81 and find a similar transit-timing variations behaviour for WASP-81b with an offset of approximately 30 s compared to the value we produce here.

<sup>4</sup> 78 expected measurements for WASP-53 and 60 for WASP-81 according to the following tool: <http://gaia.esac.esa.int/gost/>

<sup>5</sup> We do not detect a secondary set of lines in either of our spectra. This is equivalent to a limit of  $i_c > 2^\circ$  and  $>5^\circ$ , respectively.



**Figure 9.** MEGNO maps showing regions of stability in green, and chaos in red, showing where planets with the masses of WASP-53b and WASP-81b could exist. Any value in excess of 4 has the same colour. All regions on the left-hand side of the graphs are stable, and all regions on the right-hand side are unstable (within the outer orbit). Maps computed by integrating for  $5 \times 10^4$  yr.



**Figure 10.** MEGNO maps showing regions of stability in green, and chaos in red, for massless particles between the positions of the inner and outer objects in the WASP-53 and WASP-81 systems. Any value in excess of 4 has the same colour. Maps computed by integrating for  $5 \times 10^4$  yr.

### 7.3 Estimating $k_2$

Secular interactions between pairs of orbiting planets usually excite their orbital eccentricities. In the case of WASP-53 and WASP-81, the innermost planetary orbit will be affected by tidal forces that tend to damp eccentricity, while the outer, massive companion occupies a highly eccentric orbit that will excite the inner planet's eccentricity. Mardling (2007) investigated this secular problem and found that the inner planet will reach an equilibrium eccentricity, called a *fixed-point*, with a value dependent on the planet's internal density profile, and parametrized by  $k_2$ , the tidal Love number (Sterne 1939). The more mass that gets included into the core of the planet, the larger will be the fixed-point eccentricity (Batygin et al. 2009). As such a measure of the eccentricity of the tidally damped, inner orbit can yield the core-mass fraction of exoplanets. Recently, Buhler et al. (2016) have investigated the case of the HAT-P-13 system that presents an architecture similar to WASP-53 and WASP-81, and managed to constrain the core mass of the transiting hot Jupiter to  $11 M_{\oplus}$  by measuring an eccentricity of  $0.007 \pm 0.001$ . We see here how our two systems compare.

We estimated the inner eccentricity, first assuming no tidal damping, using equation (36) in Mardling (2007), and found fixed-point eccentricities of 0.00053 for WASP-53b and 0.0027 for WASP-81b (HAT-P-13's parameters yield 0.0087). Any tidal dissipation will reduce these values by an amount dependent on the internal composition. Similarly, any mutual inclination between the outer and inner orbits will reduce these values. Sadly, because the eccentricities we expect of WASP-53b and WASP-81b are so small, we find it unlikely that their internal composition will be determined soon. If however, the eccentricity is one day measured, we should expect apsidal alignment or anti-alignment between the inner and outer orbits. The Love number can also be extracted from the way a transit light curve gets affected by apsidal precession (Ragozzine & Wolf 2009).

### 7.4 Possible scenarios for the planets' formation and orbital evolution

We here speculate about the sequence of events leading to systems like WASP-53 and WASP-81, and also to HAT-P-13b (Bakos



et al. 2009), HD 59686Ab (Ortiz et al. 2016) and others. The presence of both a planet and a brown dwarf mass object (one on the planetary side and the other on the stellar side of the brown dwarf desert; Grether & Lineweaver 2006; Sahlmann et al. 2011) within the same system possibly suggests that core accretion and gravitational collapse can both operate within the same disc environment. They may also be smaller fragments from the nebula that created WASP-53A and WASP-81A.

Both WASP-53b and WASP-81b could initially have formed on circumbinary orbits, to be captured by the primary following a dynamical instability. This sort of scenario has been investigated by Sutherland & Fabrycky (2016), and shown to be an unlikely outcome, which, when it does happen, favours capture of the planet by the secondary instead of the primary. Therefore, WASP-53b and WASP-81b most likely formed within a disc surrounding WASP-53A and WASP-81A.

We can think of two alternate scenarios and make an appeal for theorists to investigate them since they could teach us much about how gas giants form.

WASP-53c and WASP-81c are too massive compared to the protoplanetary disc to have followed a type-II migration (Duffell et al. 2014; Dürmann & Kley 2015). If no or little orbital evolution has happened within the two systems that we studied, then WASP-53b and WASP-81b must have formed well within the snow line, on orbits shorter than 0.15 and 0.3 au, respectively. Theoretical work by Batygin et al. (2016) and Lee & Chiang (2015) suggests that the formation of gas giants within 1 au is feasible. According to Fung, Shi & Chiang (2014) gas can flow through the gap carved by a planet, particularly so if it reaches masses near that of a brown dwarf. The more massive the object, the more perturbations are produced at the outer gap's edge, launching streams of gas that replenish the inner disc to provide enough mass to allow the formation of a gas giant. The large eccentricities of WASP-53c and WASP-81c might have enhanced this effect. However, Lambrechts, Johansen & Morbidelli (2014) and Rosotti et al. (2016) find that planets more massive than  $20\text{--}30\text{ M}_{\oplus}$  prevent the flow of dust grains across the same gap. Accordingly, if WASP-53b and WASP-81b formed in situ via core accretion, they could only have used solids within about 0.2 au, before being able to accrete gas. If this scenario is correct, systems like WASP-53 and WASP-81 can inform us about the efficiency of core accretion, as well the minimum core mass necessary to accrete significant gas envelopes. This would also leave the planets poorer in metals than otherwise, something which is possible to determine via transmission spectroscopy (Seager & Deming 2010; Madhusudhan et al. 2014).

WASP-53c and WASP-81c are both eccentric. It has been argued that disc–planet interactions can excite the eccentricity of gap-opening planets, but not to the values that we observe for WASP-53c and WASP-81c (e.g. Goldreich & Sari 2003; D'Angelo, Lubow & Bate 2006; Teyssandier & Ogilvie 2016; Rosotti et al. 2017). This might imply that WASP-53c and WASP-81c reached their current orbital parameters after disc dispersal, possibly due to dynamical interactions with third, yet unseen companions to WASP-53 and WASP-81. If this is the case then either WASP-53b and WASP-81b disc migrated well before WASP-53c and WASP-81c reached their current orbits or WASP-53b and WASP-81b reached their current orbit following a high-eccentricity migration produced by the same instability that left their outer companions on eccentric orbital paths. In either of those cases, we expect a significant mutual inclination between the inner and outer orbits, which *Gaia* should in principle be able to measure. Coplanarity would favour the scenario outlined in the previous paragraph.

## 8 CONCLUDING WORDS

WASP-53 and WASP-81 are peculiar systems composed of both a planet and a brown dwarf. This orbital set-up is a relic of its past formation. Investigating them in further studies, notably with the help of *Gaia*, will prove invaluable for understanding planet formation and the subsequent orbital evolution, but also the relation between planet formation, brown dwarf formation and stellar formation.

*Note.* Dates are given in the BJD-TDB standard. The radii we used for Jupiter and the Sun are the volumetric mean radii. For clarity, we used the subscripts  $\star$  for the star, b for the inner planet and c for the outer object, all throughout.

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References to exoplanetary systems were obtained through the use of the paper repositories, ADS and arXiv, but also through frequent visits to the [exoplanet.eu](http://exoplanet.eu) (Schneider et al. 2011) and [exoplanets.org](http://exoplanets.org) (Wright et al. 2011) websites.

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## SUPPORTING INFORMATION

Supplementary data are available at [MNRAS](http://mnras.oxfordjournals.org/) online.

**Table A1.** Radial velocities of WASP-53 obtained with CORALIE before its recent upgrade.

**Table A2.** Radial velocities of WASP-53 obtained with CORALIE after its recent upgrade.

**Table A3.** Radial velocities of WASP-53 obtained with HARPS.

**Table A4.** Photometric observations on WASP-53 obtained with TRAPPIST.

**Table A5.** Radial velocities of WASP-81 obtained with CORALIE before its recent upgrade.

**Table A6.** Radial velocities of WASP-81 obtained with CORALIE before its recent upgrade.

**Table A7.** Radial velocities of WASP-81 obtained with HARPS.

**Table A8.** Photometric observations on WASP-81 obtained with EulerCam.

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## APPENDIX A: JOURNAL OF OBSERVATIONS

**Table A1.** Radial velocities of WASP-53 obtained with CORALIE before its recent upgrade. BJD is the Barycentric Julian Date – 245 0000 d.  $V_{\text{rad}}$  is the radial velocity obtained by fitting a cross-correlation function with a Gaussian and  $\sigma_{\text{RV}}$  is the error on  $V_{\text{rad}}$ . FWHM is the full width at half-maximum of the cross-correlation function and contrast is the amplitude. The full table is available online.

BJD (d)	$V_{\text{rad}}$ (km s <sup>-1</sup> )	$\sigma_{\text{RV}}$ (km s <sup>-1</sup> )	FWHM (km s <sup>-1</sup> )	Contrast (per cent)	Slope bisector span (km s <sup>-1</sup> )	Exposure (s)
55535.622498	– 4.38964	0.01414	7.76355	37.811	0.00572	1800.679
55563.633734	– 4.75981	0.01675	7.83813	38.764	– 0.01915	1800.736
55565.620787	– 4.10113	0.02222	7.85708	29.332	0.06569	1800.743
55583.571236	– 4.53580	0.01833	7.76400	37.735	– 0.01908	1800.754
55586.577359	– 4.45769	0.01936	7.88629	36.602	0.03945	1800.755
...	...	...	...	...	...	...

**Table A2.** Radial velocities of WASP-53 obtained with CORALIE after its recent upgrade. BJD is the Barycentric Julian Date – 245 0000 d.  $V_{\text{rad}}$  is the radial velocity obtained by fitting a cross-correlation function with a Gaussian and  $\sigma_{\text{RV}}$  is the error on  $V_{\text{rad}}$ . FWHM is the full width at half-maximum of the cross-correlation function and contrast is the amplitude. The full table is available online.

BJD (d)	$V_{\text{rad}}$ (km s <sup>-1</sup> )	$\sigma_{\text{RV}}$ (km s <sup>-1</sup> )	FWHM (km s <sup>-1</sup> )	Contrast (per cent)	Slope bisector span (km s <sup>-1</sup> )	Exposure (s)
56989.740666	– 3.92092	0.02509	7.82426	42.096	– 0.11250	1800.683
57001.640882	– 3.53647	0.02300	7.73609	42.361	0.00145	1800.780
57003.640065	– 4.11709	0.01937	7.71262	41.800	– 0.01624	1800.075
57004.682396	– 3.59763	0.02263	7.74968	41.276	– 0.00541	1800.932
57012.592977	– 3.75798	0.02566	7.80047	42.030	– 0.03335	1800.233
...	...	...	...	...	...	...

**Table A3.** Radial velocities of WASP-53 obtained with HARPS. BJD is the Barycentric Julian Date – 245 0000 d.  $V_{\text{rad}}$  is the radial velocity obtained by fitting a cross-correlation function with a Gaussian and  $\sigma_{\text{RV}}$  is the error on  $V_{\text{rad}}$ . FWHM is the full width at half-maximum of the cross-correlation function and contrast is the amplitude. The full table is available online.

BJD (d)	$V_{\text{rad}}$ (km s <sup>-1</sup> )	$\sigma_{\text{RV}}$ (km s <sup>-1</sup> )	FWHM (km s <sup>-1</sup> )	Contrast (per cent)	Slope bisector span (km s <sup>-1</sup> )	Exposure (s)
55802.804248	– 3.87536	0.01428	6.35185	48.819	0.01421	900.000
55802.902737	– 3.77124	0.00615	6.44231	49.068	0.01648	900.000
55803.735877	– 3.46797	0.00592	6.43227	49.138	0.02630	900.000
55803.876858	– 3.49415	0.01413	6.44454	48.445	0.03957	900.000
55825.652716	– 4.00995	0.00918	6.43519	50.675	– 0.00062	900.006
...	...	...	...	...	...	...

**Table A4.** Photometric observations on WASP-53 obtained with TRAPPIST. HJD is the Heliocentric Julian Date – 245 0000 d.  $F$  is the normalized flux and  $\sigma_F$  is the error on  $F$ .  $\Delta x$  and  $\Delta y$  are the offsets in the position of the PSF,  $z$  is the airmass and FWHM is the full width at half-maximum of the PSF.  $F_{\text{sky}}$  is the sky brightness, and finally we present the exposure time. The full table is available online. Data for other transits are presented in a similar manner.

HJD <sub>UTC</sub> (d)	$F$	$\sigma_F$	$\Delta x$ (pixel)	$\Delta y$ (pixel)	$z$	FWHM (pixel)	$F_{\text{sky}}$ (e <sup>−</sup> pixel <sup>−1</sup> )	Exposure (s)
55817.69393	1.00258	0.000292	0.00	0.00	1.29	5.4816	833.668	12.0
55817.69419	1.00248	0.000292	−1.75	−5.59	1.29	4.9724	835.230	12.0
55817.69441	0.99973	0.000291	−1.81	−5.85	1.28	4.8800	836.924	12.0
55817.69462	0.99925	0.000291	−2.23	−5.80	1.28	4.7187	837.804	12.0
55817.69485	0.99821	0.000290	−2.49	−5.87	1.28	4.6237	838.849	12.0
...	...	...	...	...	...	...	...	...

**Table A5.** Radial velocities of WASP-81 obtained with CORALIE before its recent upgrade. BJD is the Barycentric Julian Date – 245 0000 d.  $V_{\text{rad}}$  is the radial velocity obtained by fitting a cross-correlation function with a Gaussian and  $\sigma_{\text{RV}}$  is the error on  $V_{\text{rad}}$ . FWHM is the full width at half-maximum of the cross-correlation function and contrast is the amplitude. The full table is available online.

BJD (d)	$V_{\text{rad}}$ (km s <sup>−1</sup> )	$\sigma_{\text{RV}}$ (km s <sup>−1</sup> )	FWHM (km s <sup>−1</sup> )	Contrast (per cent)	Slope bisector span (km s <sup>−1</sup> )	Exposure (s)
55833.506244	−60.15011	0.02609	7.87708	28.619	−0.02494	1800.760
55834.511338	−60.31782	0.01861	8.06275	28.959	0.00051	1800.778
55835.513064	−60.23827	0.02606	7.95214	29.089	−0.10707	1800.742
55851.553175	−60.35013	0.02264	7.95814	26.734	−0.00195	1800.720
55852.509304	−60.17890	0.02674	7.99858	27.682	0.01749	1800.738
...	...	...	...	...	...	...

**Table A6.** Radial velocities of WASP-81 obtained with CORALIE before its recent upgrade. BJD is the Barycentric Julian Date – 245 0000 d.  $V_{\text{rad}}$  is the radial velocity obtained by fitting a cross-correlation function with a Gaussian and  $\sigma_{\text{RV}}$  is the error on  $V_{\text{rad}}$ . FWHM is the full width at half-maximum of the cross-correlation function and contrast is the amplitude. The full table is available online.

BJD (d)	$V_{\text{rad}}$ (km s <sup>−1</sup> )	$\sigma_{\text{RV}}$ (km s <sup>−1</sup> )	FWHM (km s <sup>−1</sup> )	Contrast (per cent)	Slope bisector span (km s <sup>−1</sup> )	Exposure (s)
57186.791639	−60.24486	0.03042	7.95261	34.569	0.01499	2700.846
57194.736148	−60.23109	0.04093	7.89406	34.000	0.06709	2700.943
57211.722549	−60.44937	0.06100	7.91874	33.627	−0.06934	2700.764
57256.542816	−60.39378	0.04425	7.97087	32.551	−0.07702	2700.014
57271.550777	−60.56686	0.04207	7.93436	32.440	−0.11635	2452.822
...	...	...	...	...	...	...

**Table A7.** Radial velocities of WASP-81 obtained with HARPS. BJD is the Barycentric Julian Date – 245 0000 d.  $V_{\text{rad}}$  is the radial velocity obtained by fitting a cross-correlation function with a Gaussian and  $\sigma_{\text{RV}}$  is the error on  $V_{\text{rad}}$ . FWHM is the full width at half-maximum of the cross-correlation function and contrast is the amplitude. One datum, which was not used in the analysis, is highlighted with an asterisk. The full table is available online.

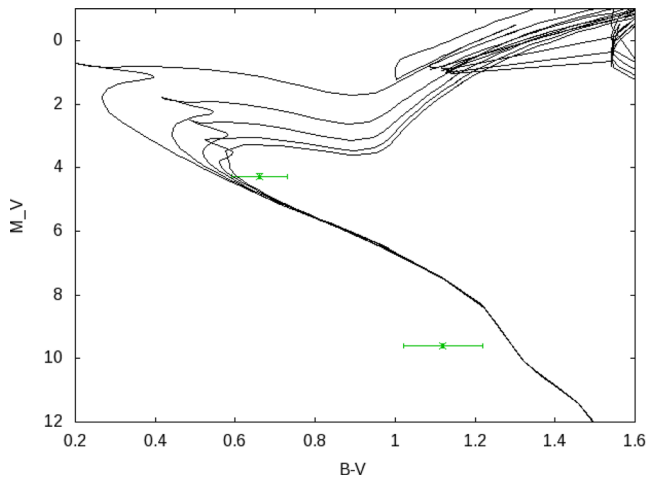
BJD (d)	$V_{\text{rad}}$ (km s <sup>−1</sup> )	$\sigma_{\text{RV}}$ (km s <sup>−1</sup> )	FWHM (km s <sup>−1</sup> )	Contrast (per cent)	Slope bisector span (km s <sup>−1</sup> )	Exposure (s)
56403.902377	−60.75558	0.01422	6.77096	36.778	0.00282	600.000
56407.886489	−60.90138	0.01261	6.79512	36.264	0.02080	600.001
56411.902064*	−73.15969	0.09454	1.71111	2.286	789.3231	600.000
56438.878243	−60.76828	0.01674	6.81762	37.631	−0.00578	900.000
56454.847018	−60.87117	0.01534	6.72971	36.706	−0.03163	600.001
...	...	...	...	...	...	...



**Table A8.** Photometric observations on WASP-81 obtained with EulerCam. HJD is the Heliocentric Julian Date – 245 0000 d.  $F$  is the normalized flux and  $\sigma_F$  is the error on  $F$ .  $\Delta x$  and  $\Delta y$  are the offsets in the position of the PSF,  $z$  is the airmass and FWHM is the full width at half-maximum of the PSF.  $F_{\text{sky}}$  is the sky brightness, and finally we present the exposure time. The full table is available online. Data for other transits are presented in a similar manner.

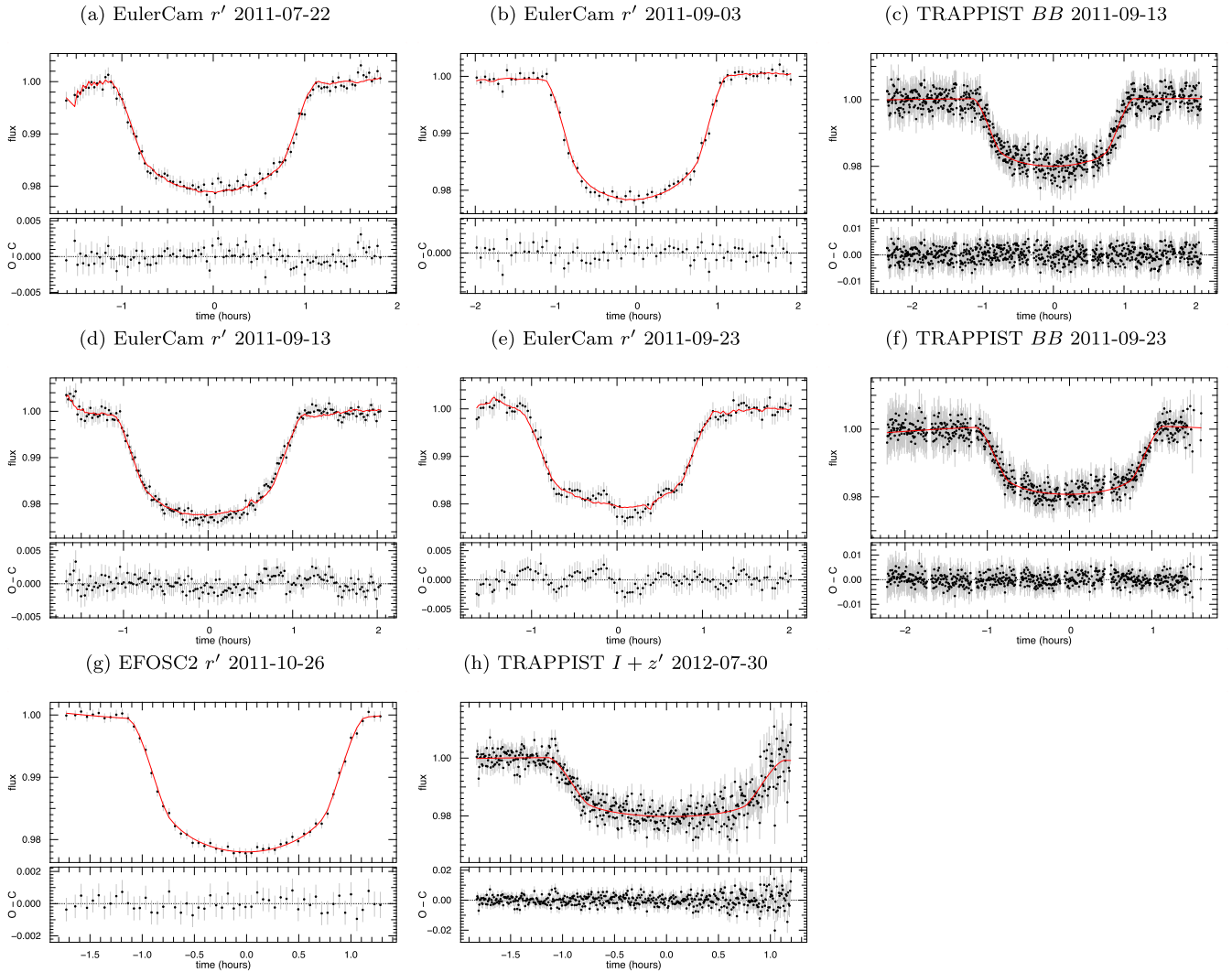
HJD <sub>UTC</sub> (d)	$F$	$\sigma_F$	$\Delta x$ (pixel)	$\Delta y$ (pixel)	$z$	FWHM (pixel)	$F_{\text{sky}}$ (e <sup>−</sup> pixel <sup>−1</sup> )	Exposure (s)
56116.673675	0.998913	0.000783	−2.21	6.16	1.34	8.77	172.69	90.0
56116.674967	0.999471	0.000784	−2.19	5.01	1.34	9.61	171.32	90.0
56116.676429	0.999586	0.000674	−1.82	6.92	1.33	8.80	228.37	120.0
56116.678103	0.998850	0.000669	−1.38	5.48	1.32	6.84	228.31	120.0
56116.679751	0.999194	0.000668	−1.67	6.92	1.32	6.69	228.53	120.0
...	...	...	...	...	...	...	...	...

## APPENDIX B: THE VISUAL COMPANION TO WASP-81

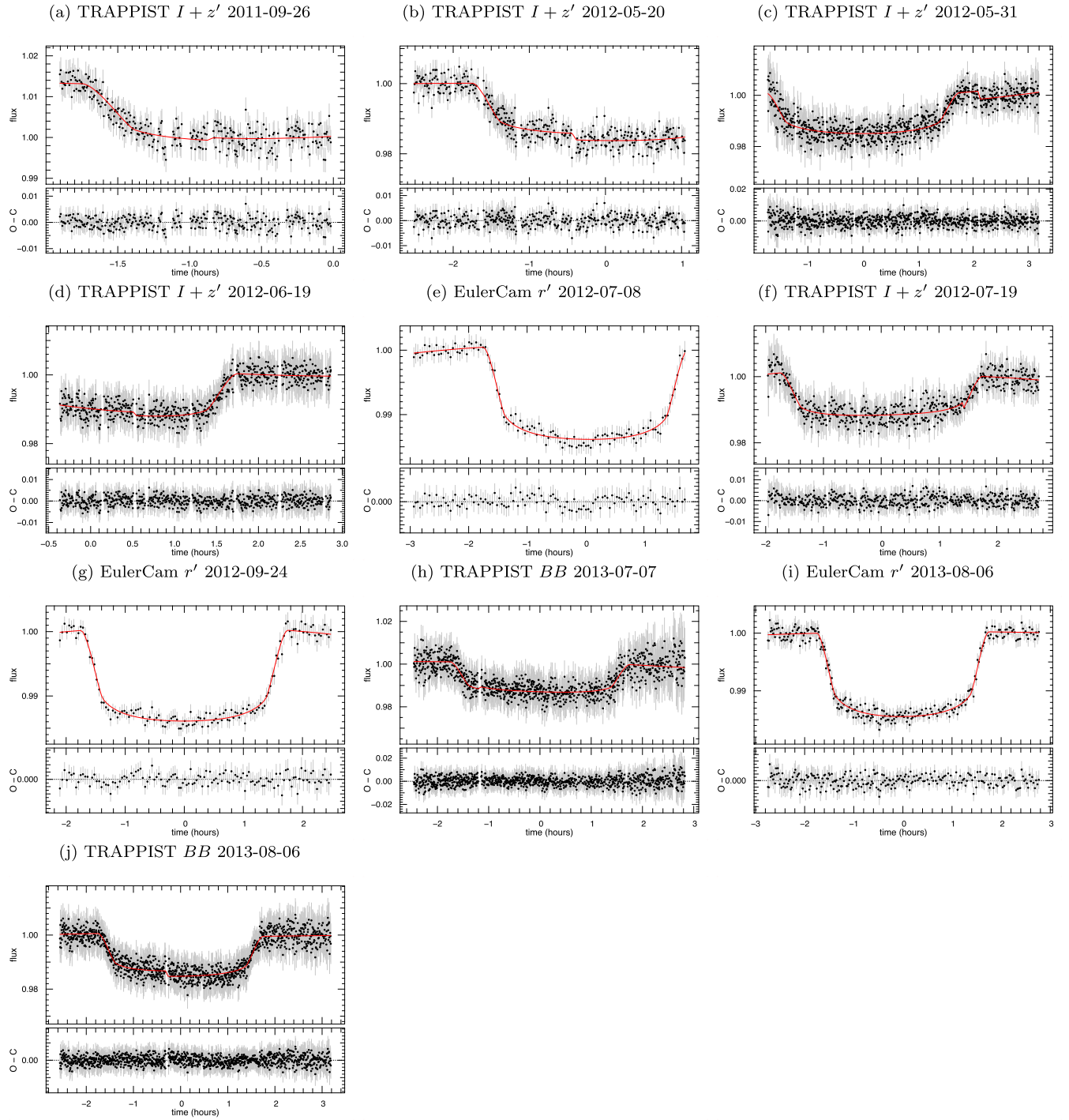


**Figure B1.** Hertzsprung–Russell diagram showing the relative positions of WASP-81A and WASP-81B assuming a similar distance. The pair is unlikely to be related. Models are from Marigo et al. (2008).

# APPENDIX C: MODELS APPLIED TO THE PHOTOMETRIC DATA



**Figure C1.** Flux as a function of time, centred around mid-transit time of WASP-53b. The red line shows the full model, including the detrending.



**Figure C2.** Flux as a function of time, centred around mid-transit time of WASP-81b. The red line shows the full model, including the detrending.

**Table C1.** Photometric time series used in this work. For each light curve this table shows the date of acquisition, the instrument and filter used, the exposure time  $T_{\text{exp}}$ , the number of data points, the baseline function selected for our global analysis (see Section 5) and the error correction factor CF used in our global analysis. For the baseline function,  $p(\epsilon^N)$  denotes, respectively, an  $N$ -order polynomial function of time ( $\epsilon = t$ ),  $x$  and  $y$  positions ( $\epsilon = xy$ );  $o$  denotes an offset at the time of a meridian flip of TRAPPIST (see Gillon et al. 2012). On two instances we also fit a sinusoidal baseline of the form  $\sin(Pt + T_0)$ , where  $P$  is the period and  $T_0$  is the phase.

System	Date	Instrument	Filter	$T_{\text{exp}}$	$N_p$	Baseline function	CF
WASP-53	2011-07-22	EulerCAM	Gunn $r'$	90 s	110	$p(t^2 + xy)$	1.3
WASP-53	2011-09-03	EulerCAM	Gunn $r'$	120 s	84	$p(t^2 + xy)$	1.8
WASP-53	2011-09-13	TRAPPIST	Blue blocking	12 s	744	$p(t^2) + o$	1.5
WASP-53	2011-09-13	EulerCAM	Gunn $r'$	80 s	127	$p(t^2 + xy + \sin(P_3t + T_{0,3}))$	1.7
WASP-53	2011-09-23	EulerCAM	Gunn $r'$	180 s	77	$p(t^2 + xy + \sin(P_4t + T_{0,4}))$	2.3
WASP-53	2011-10-26	EFOSC2	Gunn $r'$	150 s	54	$p(t^2)$	1.0
WASP-53	2012-07-30	TRAPPIST	Blue blocking	12 s	603	$p(t^2)$	1.5
WASP-53	2012-11-03	TRAPPIST	$I + z'$	15 s	464	$p(t^2)$	1.0
WASP-81	2011-09-26	TRAPPIST	$I + z'$	12 s	287	$p(t^2) + o$	0.9
WASP-81	2012-05-20	TRAPPIST	$I + z'$	25 s	372	$p(t^2) + o$	1.1
WASP-81	2012-05-31	TRAPPIST	$I + z'$	12 s	812	$p(t^2) + o$	1.4
WASP-81	2012-06-19	TRAPPIST	Blue blocking	8 s	644	$p(t^2) + o$	1.5
WASP-81	2012-07-08	EulerCAM	Gunn $r'$	120 s	116	$p(t^2)$	1.9
WASP-81	2012-07-19	TRAPPIST	$I + z'$	20 s	501	$p(t^2) + o$	1.6
WASP-81	2012-09-24	EulerCAM	Gunn $r'$	120 s	132	$p(t^2)$	1.4
WASP-81	2013-07-07	TRAPPIST	Blue blocking	10 s	917	$p(t^2) + o$	1.9
WASP-81	2013-08-06	EulerCAM	Gunn $r'$	80 s	248	$p(t^2)$	1.4
WASP-81	2013-08-06	TRAPPIST	Blue blocking	10 s	1082	$p(t^2) + o$	1.3

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